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# On the Example of Almost Pseudo-Z-symmetric Manifolds\*

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## Abstract

In the present paper we have obtained a new example of non-Ricci-flat almost pseudo-Z-symmetric manifolds in the class of equidistant spaces, which admit non-trivial geodesic mappings.

**Key words:** (pseudo-) Riemannian manifold, almost pseudo-Z-symmetric spaces, equidistant spaces.

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## 1 Introduction

In [4] was introduced an *almost pseudo-Z-symmetric space*, which is an  $n$ -dimension (pseudo-) Riemannian space  $V_n$  where the special tensor

$$Z_{ij} = R_{ij} + \varphi g_{ij},$$

satisfied the recurrent condition

$$Z_{ij,k} = (a_k + b_k)Z_{ij} + a_j Z_{ik} + a_i Z_{jk} \quad (1)$$

$R_{ij}$ ,  $g_{ij}$  and  $\varphi$  being Ricci tensor, metric tensor and scalar function.

These manifolds are generalization of symmetric and recurrent spaces which were introduced by É. Cartan [2], and A. G. Walker [19], respectively.

These manifolds were generalized in many directions, see, for example [13, pp. 292–295, 335, 338], [18]. Geodesic and holomorphically projective mappings

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of mentioned manifolds were studied in many papers too, see [6, 8, 11, 12, 13, 15, 17]. Among others, J. Mikeš [9] proved that non-Einstein Ricci-symmetric (pseudo-) Riemannian spaces ( $R_{ij,k} = 0$ ) do not admit non-trivial geodesic mappings. In paper [10] were constructed projective symmetric space which is not symmetric. For example, generalized recurrent spaces were studied in [5, 7, 14, 16].

In the paper [4] was studied almost pseudo-Z-symmetric space. As we can see, the Example 8, on p. 39–40, is false for explicit calculation. In this paper, we construct new example of these manifolds.

## 2 Equidistant manifolds

Having found the example of almost pseudo-Z-symmetric manifolds faulty [4], the present authors have constructed an example in the class of special equidistant space.

In an *equidistant space* with non isotropic concircular vector field there exists canonical coordinate system, where the metric tensor has the following form [17, pp. 92–95], [13, p. 150]:

$$ds^2 = e dx^{1^2} + f(x^1) d\tilde{s}^2, \quad (2)$$

where  $e = \pm 1$ ,  $f$  is a differentiable function and

$$d\tilde{s}^2 = \tilde{g}_{ab}(x^2, \dots, x^n) dx^a dx^b$$

is a metric of  $(n-1)$ -dimensional (pseudo-) Riemannian manifold  $\tilde{V}_{n-1}$ .

Here and after indices  $a, b, \dots = 2, 3, \dots, n$ .

In 1954 N. S. Sinyukov (see [17], [13, pp. 140-155]), thanks to their geometrical properties, gave them the name *equidistant space*. Around the year 1920 the H. W. Brinkmann [1] started studying these space and in the 1940 K. Yano [20] studied concircular vector fields. Many newly obtained results are possible to see in [3].

We denote that if  $f' \neq 0$ , then this manifold admits non-trivial geodesic mappings, see [17, 11, 13]. In the coordinate system (2) the components of metric and inverse metric tensors have the following form:

$$\begin{aligned} g_{11} &= e; & g_{1a} &= 0; & g_{ab} &= f(x^1) \tilde{g}_{ab} \\ g^{11} &= e; & g^{1a} &= 0; & g^{ab} &= f(x^1)^{-1} \tilde{g}^{ab}, \end{aligned} \quad (3)$$

where  $f (\neq 0)$  is a function of variable  $x^1$  and  $\tilde{g}_{ab}$  and  $\tilde{g}^{ab}$  are components of metric and inverse metric tensors of  $(n-1)$ -dimension on (pseudo-) Riemannian space  $\tilde{V}_{n-1}$ , their component are functions of variables  $x^2, x^3, \dots, x^n$ .

Now, non-zero components of Christoffel symbols:

$$\Gamma_{ij}^h = \Gamma_{ijk} g^{kh} \quad \text{and} \quad \Gamma_{ijk} = \frac{1}{2} (\partial_i g_{jk} + \partial_j g_{ik} - \partial_k g_{ij})$$

where  $\partial_i \equiv \partial/\partial x^i$ , have the following form:

$$\Gamma_{1ab} \equiv \Gamma_{a1b} = \frac{1}{2} f' \tilde{g}_{ab}; \quad \Gamma_{ab1} = -\frac{1}{2} f' \tilde{g}_{ab}; \quad \Gamma_{abc} = f \tilde{\Gamma}_{abc}$$

and non-zero components of Christoffel symbols of second kind:

$$\Gamma_{ab}^1 = -\frac{e}{2} f' \tilde{g}_{ab}; \quad \Gamma_{1b}^c \equiv \Gamma_{b1}^c = \frac{1}{2} \frac{f'}{f} \delta_b^c; \quad \Gamma_{ab}^c = \tilde{\Gamma}_{ab}^c \quad (4)$$

Following computation of non-zero components of the Riemannian tensor

$$R_{ijk}^h = \partial_j \Gamma_{ik}^h - \partial_k \Gamma_{ij}^h + \Gamma_{ik}^\alpha \Gamma_{\alpha j}^h - \Gamma_{ij}^\alpha \Gamma_{\alpha k}^h \quad (5)$$

$$R_{a1b}^1 \equiv -R_{ab1}^1 = -\frac{e}{2} (f'' - \frac{f'^2}{2f}) \tilde{g}_{ab},$$

$$R_{1b1}^d = -\frac{1}{2f} (f'' - \frac{f'^2}{2f}) \delta_b^d \tilde{g}_{db},$$

$$R_{abc}^d = \tilde{R}_{abc}^d - \frac{e}{4} \frac{f'^2}{f} (\tilde{g}_{ac} \delta_b^d - \tilde{g}_{ab} \delta_c^d).$$

Contracting Riemannian tensor by metric tensor, we lower indices and obtain Riemannian tensor of type  $\binom{0}{4}$

$$R_{hijk} = g_{h\alpha} R_{\alpha ijk}^{\alpha}. \quad (6)$$

After computation, we get the following non-zero components:

$$R_{1a1b} = -R_{a11b} = R_{a1b1} = R_{a11b} = -\frac{1}{2} (f'' - \frac{f'^2}{2f}) \tilde{g}_{ab}$$

$$R_{abcd} = f \tilde{R}_{abcd} - \frac{e}{4} f'^2 (\tilde{g}_{ac} \tilde{g}_{bd} - \tilde{g}_{ad} \tilde{g}_{bc}).$$

The Ricci tensor  $R_{ij} = R_{i\alpha j}^\alpha$  has these non-zero components:

$$R_{11} = R_{1\alpha 1}^\alpha = -\frac{1}{2f} (n-1) (f'' - \frac{f'^2}{2f})$$

$$R_{ab} = \tilde{R}_{ab} - \frac{e}{2} (f'' - \frac{f'^2}{2f}) \tilde{g}_{ab}.$$

### 3 Special equidistant almost pseudo-Z-symmetric spaces

The above mentioned almost pseudo-Z-symmetric spaces are defined in formula (1). Next, we shall study these spaces supposing that this space  $V_n$  is equidistant, and moreover  $\tilde{V}_{n-1}$  is Ricci flat space and component  $Z_{11}$  of tensor  $Z$  is equal to zero.

Firstly, we compute non-zero components of tensor  $Z_{ij} = R_{ij} + \varphi g_{ij}$ :

$$Z_{11} = R_{11} + \varphi(x^1) g_{11} = -\frac{1}{2f} (n-1) (f'' - \frac{f'^2}{2f}) + e\varphi;$$

$$Z_{ab} = -(\frac{e}{2} (f'' - \frac{f'^2}{2f}) - \varphi f) \tilde{g}_{ab}.$$

From our proposition ( $Z_{11} = 0$ ) it follows that the function  $\varphi$  has the following form:

$$\varphi = \frac{\epsilon}{2}(n-1)\left(f'' - \frac{f'^2}{2f}\right), \quad (7)$$

and thus

$$Z = -\frac{\epsilon n}{2}\left(f'' - \frac{f'^2}{2f}\right). \quad (8)$$

Secondly, we remember that covariant derivations of  $Z_{ij}$  have the following definition

$$Z_{ij,k} = \partial_k Z_{ij} - Z_{\alpha j} \Gamma_{ik}^\alpha - Z_{i\alpha} \Gamma_{jk}^\alpha,$$

and equation (1):

$$Z_{ij,k} = (a_k + b_k)Z_{ij} + a_j Z_{ik} + a_i Z_{jk}$$

will have the form

$$\begin{aligned} Z_{11,1} &\equiv \partial_1 Z_{11} = (3a_1 + b_1)Z_{11}; \\ Z_{11,c} &\equiv 0 = (a_c + b_c)Z_{11}; \\ Z_{1b,1} &\equiv 0 = a_b Z_{11}; \\ Z_{1b,c} &\equiv -\frac{f'}{2f} Z_{bc} + \frac{\epsilon}{2} f' Z_{11} \tilde{g}_{bc} = a_1 Z_{bc}; \\ Z_{ab,1} &\equiv \partial_1 Z_{ab} - \frac{f'}{f} Z_{ab} = (a_1 + b_1)Z_{ab}; \\ Z_{ab,c} &\equiv 0 = (a_c + b_c)Z_{ab} + a_a Z_{bc} + a_b Z_{ac}. \end{aligned}$$

Because  $Z_{11} = 0$ , the above equations are simplify to the following form:

$$-\frac{f'}{2f} Z_{bc} = a_1 Z_{bc}; \quad (9)$$

$$\partial_1 Z_{ab} - \frac{f'}{f} Z_{ab} = (a_1 + b_1)Z_{ab}; \quad (10)$$

$$(a_c + b_c)Z_{ab} + a_a Z_{bc} + a_b Z_{ac} = 0. \quad (11)$$

Naturally  $Z_{ij} \neq 0$ , then  $Z$  must not be equal to zero. Then for  $n \geq 4$  and from (11) it implies  $a_a = b_a = 0$ . From (9) and (10) follows:

$$a_1 = -\frac{1}{2} \frac{f'}{f}, \quad \text{and} \quad b_1 = -a_1 - \frac{f'}{f} \partial_1 \ln |Z|.$$

On the base of above discussion, we can formulate this theorem:

**Theorem 1** *The equidistant space with metric (2) where metric  $d\tilde{s}^2$  defined Ricci-flat space is almost pseudo-Z-symmetric space for any non-zero function  $f(x^1) \in C^3$ ,  $f'' - \frac{f'^2}{2f} \neq 0$ .*

In this space we have tensor  $Z_{ij} = R_{ij} - \varphi g_{ij}$ , where

$$\varphi = \frac{e(n-1)}{2} \left( f'' - \frac{f'^2}{2f} \right),$$

and

$$a_i = -\delta_i^1 \left( \frac{f'}{2f} \right) \quad \text{and} \quad b_i = -\delta_i^1 \frac{f'}{2f} \left( 1 - 2 \left( \ln \left| f'' - \frac{f'^2}{2f} \right| \right)' \right).$$

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