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A PROPERTY OF CONNECTIONS OF MECHANICAL SYSTEMS OF HIGHER ORDER

ANTON DEKRÉT

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ABSTRACT. In the case of regular Lagrangian L of order 1 on the tangent bundle TM of a smooth manifold M it is known that its mechanical system S_L satisfying the fundamental equation of the Lagrangian formalism $i_S d\omega_L = -dE$ determines a connection Γ which is Lagrangian, i.e. $d\omega_L(X, Y) = 0$ for all Γ -horizontal vectors X, Y on TM . In this paper the form of this property is studied in the case of higher order tangent bundles.

Let $T^r M$ be the space of all r -jets from \mathbb{R} into a smooth manifold M with source $0 \in \mathbb{R}$. Let L be a real smooth function on $T^r M$. Roughly speaking, the main idea of the Lagrangian formalism in the classical mechanics of order r consists in a construction of a symplectic structure $(T^{2r-1}M, d\omega_L)$ and a mechanical system S_L of L on $T^{2r-1}M$, the integral curves of which satisfy the Euler differential equation of order $2r$, see [6], if $r = 1$ and [8] in the case of higher order.

In [3] we constructed connections which are determined by a semispray S on $T^r M$, see also [1]. One of them given by the Lie derivation $L_S J_1$ of the canonical morphism J_1 on $T^r M$ with respect to S was studied by a number of authors first of all in the case when S is a mechanical system of a regular Lagrangian L on TM , [2, 4, 7]. For example, [4] showed that the connection $\Gamma = L_S J_1$ is Lagrangian, i.e. $\Gamma = \text{Orth}_{d\omega_L} \Gamma$, i.e. $d\omega_L(X, Y) = 0$ for all Γ -horizontal vectors X, Y on TM . In this paper we specify this property in the case of higher order. All manifolds and mappings will be smooth.

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Introduction

Throughout this paper $Tf: TM \rightarrow TN$ denotes the differential of a map $f: M \rightarrow N$. We recall some canonical properties of the manifold $T^r M$. Let (x^i) be a chart on M . Then

$$h = j_0^r(\gamma: \mathbb{R} \rightarrow M) = \left(x_0^i = \gamma^i(0), x_1^i = \frac{d\gamma^i(0)}{dt}, \dots, x_r^i = \frac{d^r \gamma^i(0)}{dt^r} \right)$$

determines the induced chart on $T^r M$. Let $\pi_k^r: T^r M \rightarrow T^k M$,

$$(x_0^i, \dots, x_k^i, \dots, x_r^i) \mapsto (x_0^i, \dots, x_k^i)$$

denote the canonical projection of r -jets onto their k -subjets. Then (π_k^r) or $V\pi_k^r$ is the abbreviated notation for the fibre manifold $\pi_k^r: T^r M \rightarrow T^k M$ or for the vector bundle of all π_k^r -vertical vectors on (π_k^r) , respectively.

There are the canonical vector fibre morphisms

$$J_1 = \sum_{p=1}^r p dx_{p-1}^i \otimes \partial/\partial x_p^i, \quad J_k = \frac{1}{k!} J_1^k = \sum_{u=k}^r \binom{u}{k} dx_{u-k}^i \otimes \partial/\partial x_u^i$$

and the canonical vector fields

$$C_1 = \sum_{p=1}^r p x_p^i \partial/\partial x_p^i, \quad C_k = \frac{1}{k!} J_1^{k-1} C_1 = \sum_{u=k}^r \binom{u}{k} x_{u-k+1}^i \partial/\partial x_u^i$$

on $T^r M$, where $k = 2, \dots, r$. Readers are referred for constructions of these objects to [8] or [3].

Recall the following embedding $i_r: T^{r+1} M \rightarrow TT^r M$,

$$\begin{aligned} h &= j_0^{r+1} \gamma(t) \mapsto j_0^1(t \mapsto j_{s=0}^r \gamma(t+s)), \\ (x_0^i, \dots, x_{r+1}^i) &\mapsto (x_0^i, \dots, x_r^i, x_1^i, \dots, x_{r+1}^i). \end{aligned}$$

By [8], a semispray on $T^r M$ (a differential equation of order $r+1$), is a vector field S on $T^r M$ such that $J_1 S = C_1$, i.e. S is of the expression

$$S = \sum_{j=0}^{r-1} x_{j+1}^i \partial/\partial x_j^i + b^i(x_0, \dots, x_r) \partial/\partial x_r^i.$$

It is clear that $S: T^r M \rightarrow TT^r M$ determines a unique section $\bar{S}: T^r M \rightarrow T^{r+1} M$ such that $S = i_r \cdot \bar{S}$, $\bar{S}(x_0^i, \dots, x_r^i) = (x_0^i, \dots, x_r^i, b^i)$. Then $S \mapsto \bar{S}$ is

a bijection between the set of all semisprays on $T^r M$ and the set of all sections of the fibred manifold (π_r^{r+1}) .

To construct connections from a semispray S we introduced in [3] a map $\tau_S: V\pi_0^r \rightarrow TT^r M$ as follows. Let J_1 be the canonical vector bundle morphism on $T^{r+1}M$ and $h \in T^{r+1}M$, $u = \pi_r^{r+1}h$. Let $J_1^h: T_u T^r M \rightarrow (V\pi_0^{r+1})_h$ denote a vector morphism such that $J_1^h(Y) = J_1(Z)$, where $Z \in T_h T^{r+1}M$ and $T\pi_r^{r+1}(Z) = Y$. Then we define a vector bundle morphism $\bar{J}_1: V\pi_0^{r+1} \rightarrow TT^r M$ such that for $W \in (V\pi_0^{r+1})_h$ there holds $\bar{J}_1(W) = Y \in T_u T^r M$, where $J_1^h(Y) = W$. It is clear that

$$\bar{J}_1(x_0^i, \dots, x_{r+1}^i, 0, c_1^i, \dots, c_{r+1}^i) = (x_0^i, \dots, x_r^i, c_1^i, \frac{1}{2}c_2^i, \dots, \frac{1}{r+1}c_{r+1}^i).$$

Let $\bar{S}: T^r M \rightarrow T^{r+1}M$ be a semispray. Using the restriction of $T\bar{S}$ on $V\pi_0^r$ we set

$$\tau_S = \bar{J}_1 \cdot T\bar{S}|_{V\pi_0^r} = \sum_{s=1}^r \frac{1}{s} dx_s^i \otimes \partial/\partial x_{s-1}^i + \frac{1}{r+1} \sum_{p=1}^r \frac{\partial b^i}{\partial x_p^i} dx_p^j \otimes \partial/\partial x_r^i.$$

Recall some needed facts about connections on a fibred manifold $\pi: Y \rightarrow M$. A 1-form ω on Y is said to be π -semibasic if $\omega(X) = 0$ for any $X \in V\pi$. A connection Γ on Y is determined by its horizontal projection h_Γ that is a π -semibasic tangent value 1-form on Y such that $T\pi \cdot h(X) = T\pi(X)$. In a chart (x^i, y^α) on Y $h_\Gamma = dx^i \otimes \partial/\partial x^i + \Gamma_i^\alpha dx^i \otimes \partial/\partial y^\alpha$. Then $v_\Gamma = \text{id}_{TY} - h_\Gamma = (dy^\alpha - \Gamma_i^\alpha dx^i) \otimes \partial/\partial y^\alpha$ is the vertical projection of Γ and $H\Gamma$ denotes the vector fibred manifold of all Γ -horizontal vectors $X \in \text{Im } h = \text{Ker } v_\Gamma$. It is obvious that if h_Γ is the horizontal projection of a given connection Γ on Y and $\varphi: Y \rightarrow T^*M \otimes VY$ is a π -semibasic $V\pi$ -value 1-form on Y , then $h_\Gamma + \varphi$ is the horizontal projection of the other connection on Y denoted by $\Gamma + \varphi$.

Let J_k or S be the canonical vector bundle morphism or a semispray on $T^r M$, respectively. In [3] we showed that $k! \tau_S^k \cdot J_k$ is the horizontal projection of a connection ${}^{r-k}\Gamma_S$ on the fibre manifold (π_{r-k}^r) . We are interested in the connections ${}^{r-1}\Gamma_S$ and ${}^0\Gamma_S$. Their coordinate expressions are as follows

$$\begin{aligned} {}^{r-1}\Gamma_S: \tau_S J_1 &= dx_0^i \otimes \partial/\partial x_0^i + \dots + dx_{r-1}^i \otimes \partial/\partial x_{r-1}^i \\ &+ \frac{1}{r+1} \sum_{p=1}^r \frac{\partial b^i}{\partial x_p^i} p dx_{p-1}^j \otimes \partial/\partial x_r^i, \end{aligned} \quad (1)$$

$${}^0\Gamma_S: r! \tau_S^r J_r = dx_0^i \otimes \partial/\partial x_0^i + r! \sum_{s=1}^r B_{r,s}^i \partial/\partial x_s^i, \quad (2)$$

where

$$B_{n,s}^i = \frac{1}{r-n+s+1} \cdots \frac{1}{r+1} A_s^i, \quad A_1^i = \frac{\partial b^i}{\partial x_r^j} dx_0^j$$

and for $s = 2, \dots, r$

$$A_s^i = \frac{\partial b^i}{\partial x_{r-s+1}^j} \frac{1}{r-s+2} \cdots \frac{1}{r} dx_0^j + \sum_{j=1}^{s-1} \frac{\partial b^i}{\partial x_{r-s+1+j}^u} B_{s-1,j}^u.$$

Remark. The connection ${}^{r-1}\Gamma_S$ coincides with the one determined by $L_S J_1$, see [8]. Other connections on (π_{r-k}^r) can be constructed from S by π_{r-k}^r -semibasic $V\pi_{r-k}^r$ -value 1-forms on $T^r M$. For example such forms are $J_k + L_S J_{k+1}, \dots, J_{r-1} + L_S J_r$. We refer to [3] for details.

Connections induced by mechanical systems of higher order

First we will give a brief survey of the main classical mechanics notions on a smooth manifold M . We refer to [8] for detailed information on higher order mechanics.

Every canonical morphism J_s on $T^r M$ determines a derivative of first order $d_{J_s} = [i_{J_s}, d] = i_{J_s} d - d i_{J_s}$, where d denotes the standard exterior derivative and

$$i_{J_s} \omega(X_1, \dots, X_p) = \sum_{j=1}^p \omega(X_1, \dots, J_s(X_j), \dots, X_p).$$

Let $f: T_1^k M \rightarrow \mathbb{R}$ be a smooth function. Regard df as a real function on $TT^k M$. Let $i_k: T^{k+1} M \rightarrow TT^k M$ be the above recalled canonical embedding. Then $d_T f = i_k^*(df) = \sum_{p=0}^k f_{x_p^i} x_{p+1}^i$ is a function on $T^{k+1} M$, where we introduced the notation $f_{x_p^i} := \frac{f}{\partial x_p^i}$ for further use. It is clear that d_T is a derivation operator which can be extended to a derivative (due to Tulczyjew, [9]) of order 0 commutative with d in the algebra λ that is the quotient set of $\bigcup_k \lambda(T^k M)$ by the equivalence relation according to which two forms $\alpha \in \lambda(T^k M)$, $\beta \in \lambda(T^j M)$, $k \geq j$, are equivalent if $\alpha = (\pi_j^k)^* \beta$, see [8].

Let $L: T^r M \rightarrow \mathbb{R}$ be a real function (Lagrangian of order r). Then the mechanical system of L is a vector field S_L on $T^{2r-1} M$ such that

$$i_{S_L} d\omega_L = -dE, \tag{3}$$

where

$$\omega_L = \sum_{i=1}^r (-1)^i d_T^{i-1} d_{J_i} L, \quad E = \sum_{i=1}^r (-1)^i d_T^{i-1} C_i(L) + L.$$

By induction it can be checked that

$$\omega_L = \sum_{p=1}^r R_i^p dx_{p-1}^i, \quad E = \sum_{p=1}^r R_i^p x_p^i + L, \quad (4)$$

where

$$R_i^p = \sum_{q=0}^{r-p} (-1)^{q+1} d_T^q L_{x_{p+q}^i}, \quad p = 1, \dots, r, \quad (5)$$

$$d_T^q L_{x_i^i} = \sum_{n=1}^q L_{x_s}^i x_{p_1}^{i_1} \dots x_{p_n}^{i_n} \sum_{t_1 + \dots + t_u = q} A_{t_1 \dots t_u} x_{p_1+t_1}^{i_1} \dots x_{p_u+t_u}^{i_u},$$

where $t_1 \leq t_2 \leq \dots \leq t_u$ are positive integers and

$$A_{(t_{11} \dots t_{1s_1}) \dots (t_{k1} \dots t_{ks_k})} = \frac{q!}{(t_{11}!)^{s_1} s_1! \dots (t_{k1}!)^{s_k} s_k!},$$

$$t_{j1} = t_{j2} = \dots = t_{j_{s_j}}.$$

Let us emphasize that R_i^p is a local function of variables x_0^i, \dots, x_{2r-p}^i .

LEMMA 1. *A semispray $S = \sum_{s=0}^{2r-2} x_{s+1}^i \partial / \partial x_s^i + b^i \partial / \partial x_{2r-1}^i$ is a solution of the equation $i_S d\omega_L = -dE$ if and only if $dR_i^1(S) = -L_{x_0^i}$.*

Proof.

$$d\omega_L = \sum_{p=1}^r dR_i^p \wedge dx_{p-1}^i, \quad dE = \sum_{p=1}^r (x_p^i dR_i^p + R_i^p dx_p^i) + \sum_{p=0}^r L_{x_p^i} dx_p^i.$$

Then the equation $i_S d\omega_L = -dE$ is of the form

$$\sum_{p=0}^{r-1} dR_i^{p+1}(S) dx_p^i + \sum_{p=1}^r R_i^p dx_p^i + \sum_{p=0}^r L_{x_p^i} dx_p^i = 0$$

that is satisfied if and only if

$$dR_i^1(S) + L_{x_0^i} = 0, \quad (6)$$

$$dR_i^{p+1}(S) + R_i^p + L_{x_p^i} = 0, \quad p = 1, \dots, r-1, \quad (7)$$

$$R_i^r + L_{x_r^i} = 0. \quad (8)$$

According to (5) the relation (8) is correct. It is easy to see that $d_T R_i^{p+1} = -R_i^p - L_{x_p^i}$. Then (7) is satisfied if and only if $dR_i^{p+1}(S) = d_T R_i^{p+1}$. This is right since

$$dR_i^{p+1} = \sum_{q=0}^{r-p-1} (-1)^{q+1} d_T^q \left(\sum_{v=1}^r L_{x_{p+q+1}^i} x_v^j dx_v^j \right).$$

Lemma 1 is proved.

The form ω_L being a π_{r-1}^{2r-1} -semibasic 1-form on $T^{2r-1}M$ determines a fibre morphism $\mathcal{L}_L: T^{2r-1}M \rightarrow T^*T^{r-1}M$ over $\text{id}_{T^{r-1}M}$ such that $\mathcal{L}_L(h)(Y) = \omega_L(X)$, where $X \in T_h T^{2r-1}M$ and $T\pi_{r-1}^{2r-1}(X) = Y$. In the induced chart $(x_0^i, \dots, x_{r-1}^i, z_i^0, \dots, z_i^{r-1})$ on $T^*T^{r-1}M$ the map \mathcal{L}_L (the Legendre transformation of L) is of the form

$$\bar{x}_{p-1}^i = x_{p-1}^i, \quad z_i^{p-1} = R_i^p(x_0^i, \dots, x_{2r-p}^i).$$

Then adding the equations

$$\begin{aligned} d\bar{x}_{p-1}^i &= dx_{p-1}^i, \\ dz_i^{p-1} &= dR_i^p = \sum_{u=0}^{2r-p-1} \frac{\partial R_i^p}{\partial x_u^j} dx_u^j + (-1)^{r-p+1} L_{x_{2r-p}^i} dx_{2r-p}^j, \end{aligned} \quad (9)$$

which follow from (5) we get the tangent prolongation $T\mathcal{L}_L$ of the Legendre transformation of L .

Denote by $p_{2r-2}^{2r-1}, \dots, p_0^{2r-1}$ the following submersions from $T^*T^{r-1}M$:

$$p_s^{2r-1}: T^*T^{r-1}M \rightarrow (V\pi_{2r-2-s}^{r-1})^*,$$

$$\begin{aligned} p_s^{2r-1}(h) &= h|_{V\pi_{2r-2-s}^{r-1}} = (x_0^i, \dots, x_{r-1}^i, z_i^{2r-1-s}, \dots, z_i^{r-1}), \\ & \quad s = r, \dots, 2r-2. \end{aligned}$$

$$p_{r-1}^{2r-1}: T^*T^{r-1}M \rightarrow T^{r-1}M,$$

$$p_{r-1}^{2r-1}(h) = (x_0^i, \dots, x_{r-1}^i).$$

$$p_k^{2r-1} = \pi_k^{r-1} \cdot p_{r-1}^{2r-1}: T^*T^{r-1}M \rightarrow T^k M, \quad k = 0, \dots, r-2.$$

The equations (9) immediately give

LEMMA 2. *Let L be a Lagrangian of order r . Then the Legendre transformation $\mathcal{L}_L: T^{2r-1}M \rightarrow T^*T^{r-1}M$ is a fibre morphism from (π_u^{2r-1}) into (p_u^{2r-1}) for $u = 0, \dots, 2r - 2$.*

Let us recall some notions of geometry on $T^*T^{r-1}M$. There is the canonical Liouville form $\lambda = \sum_{p=0}^{r-1} z_p^i dx_p^i$ on $T^*T^{r-1}M$ that is a 1-form such that $\lambda(X) = z(Tp_{r-1}^{2r-1}X)$, $X \in T_z(T^*T^{r-1}M)$. Then $d\lambda = \sum_{p=0}^{r-1} dz_p^i \wedge dx_p^i$ is the canonical symplectic form.

Let K be a given connection on (p_{2r-2}^{2r-1}) having the horizontal projection of the form

$$h_K = \sum_{p=0}^{r-1} dx_p^i \otimes \partial/\partial x_p^i + \sum_{p=1}^{r-1} dz_i^p \otimes \partial/\partial z_i^p + \left(\sum_{p=1}^{r-1} K_{pi}^j dz_j^p + \sum_{p=0}^{r-1} \bar{K}_{ij}^p dx_p^j \right) \otimes \partial/\partial z_i^0. \quad (10)$$

We are interested in the connection $\text{Orth}_{d\lambda} K$, the horizontal vectors X of which are orthogonal to all K -horizontal vectors Y on $T^*T^{r-1}M$ according to the form $d\lambda$, i.e. $d\lambda(X, Y) = 0$. It is easy to verify that the connection $\text{Orth}_{d\lambda} K$ is a connection on (p_0^{2r-1}) with the horizontal projection

$$h_{\text{Orth}_{d\lambda} K} = dx_0^i \otimes \partial/\partial x_0^i - \sum_{p=1}^{r-1} K_{pj}^i dx_0^j \otimes \partial/\partial x_p^i + \sum_{p=0}^{r-1} \bar{K}_{ji}^p dx_0^j \otimes \partial/\partial z_i^p. \quad (11)$$

Return to the fundamental equation (3) of the Lagrangian formalism. Recall that the 2-form $d\omega_L = \sum_{p=1}^r dR_i^p \wedge dx_{p-1}^i$ is symplectic if and only if the forms dR_i^p , dx_{p-1}^i , $p = 1, \dots, r$, are independent in any induced chart. Then by virtue of the relation (9) $d\omega_L$ is symplectic if and only if $\det(L_{x_i^i x_r^j}) \neq 0$. In this case the Lagrangian L is called regular.

In what follows, we shall consider the Lagrangian L to be regular. Now the equation (3) $i_X d\omega_L = -dE$ has a unique solution $X = S_L$ called the

mechanical system of the Lagrangian L . By Lemma 1 S_L is the semispray on $T^{2r-1}M$ which satisfies the equation (6), i.e.

$$b^i = (-1)^{r+1} \tilde{L}^{is} \left(\sum_{u=0}^{2r-2} \frac{\partial R_s^1}{\partial x_u^j} x_{u+1}^j + Lx_0^s \right), \quad L_{x_i^i x_i^i} \tilde{L}^{sj} = \delta_i^j. \quad (12)$$

We can find relations between the connections ${}^{2r-2}\Gamma_S$ and ${}^0\Gamma_S$ on $T^{2r-1}M$ determined by the mechanical system $S = S_L$. In this case the connection $\text{Orth}_{d\omega_L} {}^{2r-2}\Gamma_S$ which is $d\omega_L$ -orthogonal to ${}^{2r-2}\Gamma_S$ is a connection on (π_0^{2r-1}) . To specify some properties of the connections ${}^0\Gamma_S$ and $\text{Orth}_{d\omega_L} {}^{2r-2}\Gamma_S$ we use the Legendre transformation \mathcal{L}_L . By (9) it is clear that \mathcal{L}_L is a local symplectic isomorphism of the symplectic manifolds $(T^{2r-1}M, d\omega_L)$ and $(T^*T^{r-1}M, d\lambda)$.

Let K and 0H be the connections on $T^*T^{r-1}M$, which are the images of the connections ${}^{2r-2}\Gamma_S$ and ${}^0\Gamma_S$ under the Legendre transformation \mathcal{L}_L . Recall that the difference ${}^0H - {}^0\bar{H}$ of two connections on (p_0^{2r-1}) is a p_0^{2r-1} -semibasic Vp_0^{2r-1} -value form on $T^*T^{r-1}M$. We will deduce that the form ${}^0H - \text{Orth}_{d\lambda} K$ is a p_0^{2r-1} -semibasic Vp_1^{2r-1} -value 1-form on $T^*T^{r-1}M$.

Let (10) be the expression of the horizontal projection of K . Let

$$O_H = dx_0^0 \otimes \partial/\partial x_0^i + \sum_{p=1}^{r-1} H_{pj}^i dx_0^j \otimes \partial/\partial x_p^i + \sum_{p=1}^{r-1} \bar{H}_{ij}^p dx_0^j \otimes \partial/\partial z_i^p$$

be the horizontal projection of the connection 0H . Then by (11) the form ${}^0H - \text{Orth}_{d\lambda} K$ is a p_0^{2r-1} -semibasic and just Vp_1^{2r-1} -value form if $H_{1j}^i = -K_{1j}^i$ and $H_{2j}^i \neq -K_{2j}^i$. To prove these relations we will find the local functions H_{1j}^i , K_{1j}^i , H_{2j}^i , K_{2j}^i .

By (1) the horizontal projection of the connection ${}^{2r-2}\Gamma_{S_L}$ is

$$\sum_{p=0}^{2r-2} dx_p^i \otimes \partial/\partial x_p^i + \frac{1}{2r} \sum_{p=1}^{2r-1} \frac{\partial b^i}{\partial x_p^j} p dx_{p-1}^j \otimes \partial/\partial x_{2r-1}^i,$$

where the local functions b^i are defined by (12). In another way, ${}^{2r-2}\Gamma_{S_L}$ is given by the equation

$$dx_{2r-1}^i = \sum_{p=0}^{2r-2} \Gamma_j^{ip} dx_p^j, \quad \Gamma_j^{ip} = \frac{1}{2r} (p+1) \frac{\partial b^i}{\partial x_{p+1}^j}.$$

Then the equation of the connection K is

$$dz_i^0 = \sum_{s=0}^{2r-4} \frac{\partial R_i^1}{\partial x_s^j} dx_s^j + \frac{\partial R_i^1}{\partial x_{2r-3}^j} dx_{2r-3}^j + \frac{\partial R_i^1}{\partial x_{2r-2}^j} dx_{2r-2}^j + (-1)^r L_{x_i^i x_i^i} \sum_{p=0}^{2r-2} \Gamma_v^{jp} dx_p^v. \quad (13)$$

Using the equations

$$dx_{2r-3}^j = (-1)^r \tilde{L}^{jw} \left(dz_w^2 - \sum_{s=0}^{2r-4} \frac{\partial R_w^3}{\partial x_s^k} dx_s^k \right)$$

$$dx_{2r-2}^j = (-1)^{r-1} \tilde{L}^{jw} \left[dz_w^1 - \sum_{s=0}^{2r-4} \frac{R_w^2}{x_s^k} dx_s^k - \frac{\partial R_w^2}{\partial x_{2r-3}^v} (-1)^r \tilde{L}^{vt} \left(dz_t^2 - \sum_{s=0}^{2r-4} \frac{\partial R_t^3}{\partial x_s^k} dx_s^k \right) \right],$$

which follow from (9), we deduce from (13) that

$$K_{1i}^j = \left(\frac{\partial R_i^1}{\partial x_{2r-2}^q} + (-1)^r L_{x_i^i x_i^i} \Gamma_q^{u,2r-2} \right) (-1)^{r-1} \tilde{L}^{qj}$$

$$K_{2i}^j = \left(\frac{\partial R_i^1}{\partial x_{2r-3}^q} + (-1)^r L_{x_i^i x_i^i} \Gamma_q^{u,2r-3} \right) (-1)^r \tilde{L}^{qj}$$

$$+ \left(\frac{\partial R_i^1}{\partial x_{2r-2}^q} + (-1)^r L_{x_i^i x_i^i} \Gamma_q^{u,2r-2} \right) (-1)^{r-1} \tilde{L}^{qw} \left(-\frac{\partial R_w^2}{\partial x_{2r-3}^v} (-1)^r \tilde{L}^{vj} \right).$$

By (2) and (9) the horizontal projection of the connection 0H is of the form

$$dx_0^i \otimes \partial / \partial x_0^i + (2r-1)! \sum_{s=1}^{r-1} B_{2r-1,s}^i \partial / \partial x_s^i + \sum_{p=1}^{r-1} \tilde{H}_{ij}^p dx_0^j \otimes \partial / \partial z_s^p.$$

Since $B_{2r-1,1}^i = \frac{1}{(2r)!} A_1^i = \frac{1}{(2r)!} \frac{\partial b^i}{\partial x_{2r-1}^j} dx_0^j$,

$$H_{1j}^i = \frac{(2r-1)!}{(2r)!} \frac{\partial b^i}{\partial x_{2r-1}^j} = \frac{1}{2r} \frac{\partial b^i}{\partial x_{2r-1}^j}.$$

Quite analogously we deduce that

$$H_{2j}^i = \frac{1}{r} \left(\frac{\partial b^i}{\partial x_{2r-2}^j} \frac{1}{2r-1} + \frac{\partial b^i}{\partial x_{2r-1}^u} \frac{1}{2r} \frac{\partial b^u}{\partial x_{2r-1}^j} \right).$$

Using $\Gamma_j^{ip} = \frac{1}{2r}(p+1) \frac{\partial b^i}{\partial x_{p+1}^j}$, (12) and

$$R_i^p = \sum_{q=0}^{r-p} (-1)^{q+1} \sum_{u=1}^q L_{x_{p+q}^i x_{p_1}^{i_1} \dots x_{p_n}^{i_n}} \sum_{t_1+\dots+t_u=q} A_{t_1 \dots t_u} x_{p_1+1}^{i_1} \dots x_{p_u+t_u}^{i_n}$$

we obtain

$$-K_{1i}^i = + \frac{1}{2r} \tilde{L}^{jq} (L_{x_r^i x_{r-1}^q} - L_{x_{r-1}^i x_r^q} - r L_{x_r^i x_{p_1}^{i_1} x_r^q x_{p_1+1}^{i_1}}) = H_{1i}^j.$$

By a little more complicated calculation it can be shown that $H_{2j}^i \neq -K_{2i}^i$. Therefore the form ${}^0H\text{-Orth}_{d\lambda} K$ is p_0^{2r-1} -semibasic with values in Vp_1^{2r-1} . Then using the symplectic isomorphism \mathcal{L}_L we get the following proposition:

PROPOSITION. *Let L be a regular Lagrangian. Then the connections ${}^0\Gamma$ and ${}^{2r-2}\Gamma$ determined by the mechanical system S_L of L are such that the form ${}^0\Gamma\text{-Orth}_{d\omega_L} {}^{2r-2}\Gamma$ is π_0^{2r-1} -semibasic with values in $V\pi_1^{2r-1}$.*

Remarks.

1. If $r = 1$, then ${}^0\Gamma = {}^{2\cdot 1-2}\Gamma$ and the Proposition asserts that ${}^0\Gamma = \text{Orth}_{d\omega_L} {}^0\Gamma$, i.e. that the connection ${}^0\Gamma$ is Lagrangian, compare with [4, 5].

2. In [3] large families γ_S^k of natural operators ϕ of first order from the space of all semisprays S on T^nM into the space of connections on (π_k^n) , $k = 0, \dots, n-1$, have been constructed. The simplest of them, from the point of their coordinate expression view, are the connections $(n-k)! \tau_S^{n-k} J_{n-k}$. In the case $n = 3$, it is proved in [5] that there are no other first order natural operators. If S_L is the mechanical system determined by a regular Lagrangian L of order r and if $\Gamma \in \gamma_{S_L}^k$, then the connection $\text{Orth}_{d\omega_{L_S}} \Gamma$ is a connection on (π_{2r-2-k}^{2r-1}) . For $r = 2$, [5] shows that $\text{Orth}_{d\omega_{L_S}} \Gamma$ does not belong to $\gamma_{S_L}^{2r-2-k}$. Our proposition demonstrates what properties the connections Γ and $\text{Orth}_{d\omega_L} \Gamma$ only have. The operators ϕ do not determine all of the connections induced by L . Every regular Lagrangian L determines other families of connections $\text{Orth}_{d\omega_L} \Gamma$, where Γ is a connection belonging to $\gamma_{S_L}^k$, $k = 0, \dots, 2r-2$.

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