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CALCULATION OF THE FREE ENERGY IN A SIMPLE MODEL

A.K. Kwaśniewski

Abstract:

The free energy per site is calculated, for the one-dimensional periodic chain, in an arbitrary external magnetic field. This chain is a one-dimensional counterpart of the two-dimensional, three state Potts models.

The use of generalized Clifford algebras, has recently resulted in revealing new perspectives for calculation of the partition function for Potts models [1,2].

At the same time the main algebraic obstacle for straightforward generalizations of the known in the Ising case model methods - seems to be localized now. Therefore it is useful to get further experience, while calculating the partition function for the one-dimensional counterparts of both standard and planar Potts models.

These are one-dimensional periodic chains with partition functions defined as follows:

$$Z_N = \sum_{\{\mu\}} \exp\{a \sum_{i=1}^N \delta(\mu_i - \mu_{i+1}) + B \sum_{i=1}^N \operatorname{Re} \mu_i\}, \quad (1)$$

$$Z'_N = \sum_{\{\mu\}} \exp\{a \sum_{i=1}^N \operatorname{Re}(\mu_i \bar{\mu}_{i+1}) + B \sum_{i=1}^N \operatorname{Re} \mu_i\}, \quad (2)$$

where $\mu_i \in \{\omega^k\}_{k=0}^2$, $\omega = \exp(i \frac{2\pi}{3})$ and $\mu_{N+1} = \mu_1$, $a \neq 0$, $a, B \in \mathbb{R}$.

The transfer matrices L and L' are given correspondingly by:

$$L(a) = \begin{pmatrix} e^a & 1 & 1 \\ 1 & e^a & 1 \\ 1 & 1 & e^a \end{pmatrix} \begin{pmatrix} e^B & 0 & 0 \\ 0 & e^{B\operatorname{Re}\omega} & 0 \\ 0 & 0 & e^{B\operatorname{Re}\omega} \end{pmatrix}, \quad (3)$$

$$L'(a) = e^{a\operatorname{Re}\omega} L(a - \operatorname{Re}\omega). \quad (4)$$

Due to (4) it is sufficient to calculate Z_N partition function only.

We shall look for eigenvalues of $L(a)$, and for that to do let us introduce the following notation:

$$e^a = \alpha, \quad b_0^2 = e^B, \quad b^2 = e^{BRe\omega}, \quad x_0 = \alpha - \lambda b_0^{-2} \equiv \alpha - \lambda\beta.$$

$$x = \alpha - \lambda b^{-2} \equiv \alpha - \lambda\beta.$$

Note that $\beta_0 \beta^2 = 1$.

Define now for the moment the A & B matrices:

$$A = \begin{pmatrix} \alpha & 1 & 1 \\ 1 & \alpha & 1 \\ 1 & 1 & \alpha \end{pmatrix}, \quad B = \begin{pmatrix} b_0 & 0 & 0 \\ 0 & b & 0 \\ 0 & 0 & b \end{pmatrix}.$$

Then one easily sees that the spectrum of $L(a)$ is that of BAB , what amounts to looking for the roots of $\det(BAB - \lambda I)$, where

$$\det(BAB - \lambda I) = \begin{vmatrix} x_0 & 1 & 1 \\ 1 & x & 1 \\ 1 & 1 & x \end{vmatrix} = (x - 1)(xx_0 + x_0 - 2). \quad (5)$$

Hence the first eigenvalue of $L(a)$ is equal to

$$\lambda_0 = (\alpha - 1)b^2. \quad (6)$$

The other two are to be found from

$$\lambda^2 - [\alpha\beta^2 + (\alpha + 1)\beta_0\beta]\lambda + \beta[\alpha(\alpha + 1) - 2] = 0. \quad (7)$$

Denote the coefficient of λ by A_1 while $A_2 = \beta[\alpha(\alpha + 1) - 2]$. Then we see that $L(a)$ has two more eigenvalues as

$$\Delta = A_1^2 - 4A_2 = [\alpha\beta^2 - (\alpha + 1)\beta_0\beta]^2 + 8\beta > 0 \quad (8)$$

for any a and B . Thus we obtain:

$$\lambda_1 = \frac{1}{2}[\alpha\beta^2 + (\alpha + 1)\beta_0\beta + \sqrt{\Delta}], \quad (9)$$

$$\lambda_2 = \frac{1}{2}[\alpha\beta^2 + (\alpha + 1)\beta_0\beta - \sqrt{\Delta}]. \quad (10)$$

It follows from (9) that $\lambda_1 > 0$ independently of a and B . Meanwhile $\lambda_2 \geq 0$ depending on the values of a only, i.e.

$$\begin{aligned} \lambda_1 > \lambda_2 > 0 & \quad \text{for } a > 0, \text{ and} \\ \lambda_2 < 0, \lambda_1 > |\lambda_2| & \quad \text{for } a < 0. \end{aligned} \tag{11}$$

In order to prove (11) it is enough to notice that

$$b^2 \lambda_1 \lambda_2 = \alpha(\alpha + 1) - 2 \quad \text{and that } \alpha > 1 \text{ for } a > 0, \text{ while } \alpha < 1 \text{ for } a < 0.$$

The inequality $\lambda_1 > |\lambda_2|$ for any α is obvious in virtue of (9) and (10). As the next important step we prove:

LEMMA 1.

Let $B \geq 0$ and $a \neq 0$, then $\lambda_1 > |\lambda_0|$. \square

Proof:

$$\frac{\lambda_1^2}{\lambda_0^2} > \frac{\lambda_1 |\lambda_2|}{\lambda_0^2} = \beta^3 \left| \frac{\alpha + 2}{\alpha - 1} \right| > 1. \quad \square$$

Now we are in a position to extend the validity of the above Lemma - to arbitrary B .

LEMMA 2.

For any B and $a \neq 0$, $\lambda_1 > |\lambda_0|$. \square

Proof:

At first we prove that the continuous function of B : $f(B) = \lambda_1 - |\lambda_0|$ never takes the zero value. The thesis to be proved then follows from Lemma 1.

Let $a > 0$. Then $\alpha > 1$ and let $\lambda_1 - |\lambda_0| = \lambda_1 - \lambda_0 = 0$. This is equivalent to $\beta_0 = \beta$ i.e. $B=0$ and this leads to contradiction as for $B=0$, $\alpha+2=\lambda_1 \neq \lambda_0 = \lambda_2 = \alpha-1$.

Let now $a < 0$, then $\alpha < 1$ and let $\lambda_1 - |\lambda_0| = \lambda_1 + \lambda_0 = 0$. Then $\alpha x + x_0 - 2 = 0$ for $\lambda = -\lambda_0$ (see (5)). This is equivalent to $x_0 = 1$ or explicitly: $\alpha^2 + \alpha^2 \kappa - \alpha \kappa - 1 = 0$; $\kappa = \beta_0 \beta^{-1} > 0$. However, this proves the contradiction as for $0 < \alpha < 1$ $\alpha^2 - 1 + \kappa(\alpha^2 - \alpha) < 0$. \square

From what was proved it follows directly that the free energy per site - $f(a, B)$ - for the model defined by (1), reads:

$$\boxed{f(a, B) = -kT \ln \lambda_1} \tag{12}$$

The free energy thus is a continuous function of all its arguments.
The further, implicate dependence of f on temperature T is through parameters a and B which naturally incorporate the $1/kT$ factor.
The 5-state model is to be presented in a forthcoming paper.

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REFERENCES

TRUONG T.T., FUB-TKM Sept.85/25(1985) Berlin
KWASNIEWSKI A.K., J.Phys.A (in press), ITP UWr 84/621 (1984) Wrocław

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