

Ewa Krot

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∂_ψ - DIFFERENCE CALCULUS BERNOULLI-TAYLOR FORMULA

EWA KROT

ABSTRACT. In this note we derive the general ∂_ψ -difference Bernoulli-Taylor formula with the rest term of the Cauchy type.

1. BERNOULLI-TAYLOR FORMULA

In [2] O.V. Viskov presents another form of the Bernoulli-Taylor formula with the rest term of the Cauchy type. For that he uses Graves-Heisenberg-Weil (GHW) algebra generators \hat{p} and \hat{q} such that:

$$(1) \quad [\hat{p}, \hat{q}] = \hat{p}\hat{q} - \hat{q}\hat{p} = 1$$

where 1 is identity operator. Using (1) and the induction one may prove the following identity:

$$(2) \quad \hat{p}\hat{q}^n = \hat{p}^n\hat{q} + n\hat{q}^{n-1} \quad (n = 1, 2, \dots).$$

Now consider the obvious identity:

$$(3) \quad \sum_{k=0}^n (\alpha_k - \alpha_{k+1}) = \alpha_0 - \alpha_{n+1}.$$

Under the substitution

$$(4) \quad \alpha_0 = 0, \quad \alpha_k = (-1)^k \frac{\hat{q}^{k-1}\hat{p}^k}{(k-1)!}, \quad k = 1, 2, \dots$$

and using (2) one can get from (3):

$$(5) \quad \hat{p} \sum_{k=0}^n \frac{(-\hat{q})^k \hat{p}^k}{k!} = \frac{(-\hat{q})^n \hat{p}^{n+1}}{n!}$$

what is Bernoulli identity (see Viskov [2]).

Example 1.1. Let \hat{p} and \hat{q} be as below:

$$\hat{p} = D \equiv \frac{d}{dx}, \quad \hat{q} = \hat{x} - y, \quad y \in F(\mathbf{R}, \mathbf{C}).$$

where

$$\hat{x}f(x) = xf(x)$$

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 The paper is in final form and no version of it will be submitted elsewhere.

for sufficiently smooth function $f : \mathbf{F} \rightarrow \mathbf{F}$.

After substitution into Bernoulli identity and application to function f as above we get:

$$D \sum_{k=0}^n \frac{(y-x)(D^k f)(x)}{k!} = \frac{(y-x)^n}{n!} (D^{k+1} f)(x).$$

Now after integration $\int_y^x dt$ we get:

$$f(y) = \sum_{k=0}^n \frac{(y-x)^k}{k!} f^{(k)}(x) + \int_x^y \frac{(y-t)^n}{n!} f^{(n+1)}(t) dt$$

what is well known Bernoulli-Taylor formula with the rest term of the Cauchy type.

2. ∂_ψ -UMBRAL CALCULUS

Now we shall present some definitions and theorems of ∂_ψ -umbral calculus. One can find more of them in [3], [6], [7].

We shall denote by \mathbf{P} the algebra of polynomials over the field \mathbf{F} of characteristic zero. Let us consider a one parameter family \mathcal{F} of sequences. Then Ψ is called admissible if $\Psi \in \mathcal{F}$. Where:

$$\mathcal{F} = \{ \Psi : \mathbf{R} \supset [a, b]; q \in [a, b] : \Psi(q) : \mathbf{Z} \rightarrow \mathbf{F}; \Psi_0(q) = 1, \\ \Psi_n(q) \neq 0, \Psi_{-n}(q) = 0, n \in \mathbf{N} \}.$$

Now let us to introduce the Ψ -notation:

$$n_\psi = \Psi_{n-1}(q)\Psi_n^{-1}(q); \\ n_\psi! = n_\psi(n-1)_\psi \cdots 2_\psi 1_\psi = \Psi_n^{-1}(q), \\ n_\psi^k = n_\psi(n-1)_\psi \cdots (n-k+1)_\psi, \\ \binom{n}{k}_\psi = \frac{n_\psi^k}{k_\psi!},$$

$$\exp_\psi \{y\} = \sum_{k=0}^\infty \frac{y^k}{k_\psi!}$$

Definition 2.1. Let $\partial_\psi : \mathbf{P} \rightarrow \mathbf{P}$ and $\partial_\psi x^n = n_\psi x^{n-1}$; ∂_ψ -linearly extended is called the Ψ -derivative.

Definition 2.2. The \hat{x}_ψ -operator (∂_ψ -multiplication operator) is the linear map such that

$$\hat{x}_\psi x^n = \frac{n+1}{(n+1)_\psi} x^{n+1} \quad n \geq 0.$$

Note that $[\partial_\psi, \hat{x}_\psi] = 1$.

Let us to introduce Ψ -multiplication $*_\psi$ of functions as specified below:

Definition 2.3.

$$(6) \quad x *_\psi x^n = \hat{x}_\psi(x^n) = \frac{(n+1)}{(n+1)_\psi} x^{n+1} \quad n \geq 0,$$

$$(7) \quad x^n *_\psi x = \hat{x}_\psi^n(x) = \frac{(n+1)!}{(n+1)_\psi!} x^{n+1} \quad n \geq 0.$$

Therefore

$$(8) \quad x *_\psi \alpha 1 = \alpha 1 *_\psi x = \alpha *_\psi x$$

and

$$f(x) *_\psi x^n = F(\hat{x}_\psi)x^n.$$

Note 2.1. For $k \neq n$, $x^n *_\psi x^k \neq x^k *_\psi x^n$ as well as $x^n *_\psi x^k \neq x^{n+k}$ – in general i.e. for arbitrary admissible Ψ .

Definition 2.4. Let us to define $*_\psi$ -powers of x according to:

$$(9) \quad x^{n*_\psi} = x *_\psi x^{(n-1)*_\psi} = \hat{x}_\psi(x^{(n-1)*_\psi}) = x *_\psi x *_\psi \dots *_\psi x = \frac{n!}{n_\psi!} x^n; \quad n \geq 0.$$

Note 2.2. Note that

$$x^{n*_\psi} *_\psi x^{k*_\psi} = \frac{n!}{n_\psi!} x^{(n+k)*_\psi} \neq \frac{k!}{k_\psi!} x^{(n+1)*_\psi} = x^{k*_\psi} *_\psi x^{n*_\psi}$$

for $k \neq n$ and $x^{0*_\psi} = 1$.

This noncommutative Ψ -multiplication $*_\psi$ is devised so as to ensure the following observations:

Observation 2.1. Let f, g be formal series. Then:

- (a) $\partial_\psi x^{n*_\psi} = n x^{(n-1)*_\psi}; \quad n \geq 0;$
- (b) $\exp_\psi[\alpha x] = \exp \alpha \hat{x}_\psi 1;$
- (c) $\exp[\alpha x] *_\psi \exp_\psi \beta \hat{x}_\psi 1 = \exp_\psi [\alpha + \beta] \hat{x}_\psi 1;$
- (d) $\partial_\psi (x^k *_\psi x^{n*_\psi}) = (Dx^k) *_\psi x^{n*_\psi} + x^k *_\psi (\partial_\psi x^{n*_\psi});$
- (e) $\partial_\psi (f *_\psi g) = (Df) *_\psi g + f *_\psi (\partial_\psi g);$ (∂_ψ -Leibnitz rule);
- (f) $f(\hat{x}_\psi)g(\hat{x}_\psi)1 = f(x) *_\psi \tilde{g}; \quad \tilde{g}(x) = g(\hat{x}_\psi)1.$

3. Ψ -INTEGRATION

Now let us to define Ψ -integration which is a right inverse operation to Ψ -derivative, i.e.

$$\partial_\psi \circ \int d_\psi t = id.$$

Note that $\partial_\psi = \hat{n}_\psi \partial_0$ where $\hat{n}_\psi x^n = (n+1)_\psi x^n; \quad n \geq 1$ and $\partial_0 x^n = x^{n-1};$ in general $(\partial_0 f)(x) = \frac{1}{x}(f(x) - f(0)).$

Definition 3.1. We define Ψ -integral as a linear operator such that

$$(10) \quad \int x^n d_\psi x = [\hat{x} \frac{1}{\hat{n}_\psi}] x^n = \hat{x} \left(\frac{1}{(n+1)_\psi} x^n \right) = \frac{1}{(n+1)_\psi} x^{n+1}; \quad n \geq 0$$

for \hat{x} as in Example 1.1.

Note 3.1. Also note that :

$$(11) \quad \partial_\psi \circ \int_\alpha^x f(t) d_\psi t = f(x)$$

and

$$(12) \quad \int_\alpha^x (\partial_\psi f)(t) d_\psi t = f(x) - f(\alpha)$$

for every formal series f .

Observation 3.1. The following formula for integration “per partes” holds:

$$\int_\alpha^\beta (f *_\psi \partial_\psi g)(x) d_\psi x = [(f *_\psi g)(x)]_\alpha^\beta - \int_\alpha^\beta ((Df) *_\psi g)(x) d_\psi x.$$

4. ∂_ψ -BERNOULLI-TAYLOR FORMULA

Let us to return to Bernoulli identity

$$(13) \quad \hat{p} \sum_{k=0}^n \frac{(-\hat{q})^k \hat{p}^k}{k!} = \frac{(-\hat{q})^n \hat{p}^{n+1}}{n!}.$$

Now let \hat{p} and \hat{q} be as below:

$$(14) \quad \hat{p} = \partial_\psi, \quad \hat{q} = \hat{z} - \psi = \hat{x}_\psi - y, \quad y \in \mathbf{F}.$$

Note that $\partial_\psi, \hat{z}_\psi = id$. After submission into (13) we get:

$$(15) \quad \partial_\psi \sum_{k=0}^n \frac{(y - \hat{x}_\psi)^k (\partial_\psi^k f)(x)}{k!} = \frac{(y - \hat{x}_\psi)^n (\partial_\psi^{n+1} f)(x)}{n!}.$$

Using (6)–(9) one can get equivalent identity:

$$(16) \quad \partial_\psi \sum_{k=0}^n \frac{(y-x)^{k*_\psi} *_\psi (\partial_\psi^k f)(x)}{k!} = \frac{(y-x)^{n*_\psi} *_\psi (\partial_\psi^{n+1} f)(x)}{n!}.$$

After integration $\int_\alpha^x d_\psi x$ using (11),(12) it gives ∂_ψ -Bernoulli-Taylor formula of the form:

$$(17) \quad f(x) = \sum_{k=0}^n \frac{1}{k!} (x-\alpha)^{k*_\psi} *_\psi (\partial_\psi^k f)(\alpha) + R_{n+1}(x)$$

with the rest term of the Cauchy type of the form:

$$(18) \quad R_{n+1}(x) = \frac{1}{n!} \int_\alpha^x (x-t)^{n*_\psi} *_\psi (\partial_\psi^{n+1} f)(t) d_\psi t.$$

Remark 4.1. In [1] is presented special case of (17),(18) which is ∂_q -Bernoulli-Taylor formula of the form:

$$(19) \quad f(x) = \sum_{k=0}^n \frac{1}{k!} (x-\alpha)^{k*_q} *_q (\partial_q^k f)(\alpha) + R_{n+1}(x)$$

with

$$(20) \quad R_{n+1}(x) = \frac{1}{n!} \int_\alpha^x (x-t)^{n*_q} *_q (\partial_q^{n+1} f)(t) d_q t,$$

where

$$(21) \quad (\partial_q f)(t) = \frac{f(t) - f(qt)}{(1-q)t}$$

and one can get it by the choice $n_\psi = n_q = 1 + q + \dots + q^{n-1}$.

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INSTITUTE OF COMPUTER SCIENCE
BIAŁYSTOK UNIVERSITY
UL. SOSNOWA 64
PL-15-887 BIAŁYSTOK
POLAND