# Acta Universitatis Carolinae. Mathematica et Physica

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Acta Universitatis Carolinae. Mathematica et Physica, Vol. 16 (1975), No. 1, 63--66

Persistent URL: http://dml.cz/dmlcz/142361

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# Inversion Integrals for the Integral Transforms Involving the Meijer's G-Function as Kernel

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Received 27 September 1973

In the present work, two integral equations with Meijer's G-function in the kernel have been solved. The technique employed in finding the solutions of integral equations is somewhat different than the techniques used in [1], [2], [6]. Later some special cases are discussed.

#### 1. Introduction

Some inversion integrals for integral equations involving either a Chebysheb or Legendre or Gegenbauer polynomial in the kernel are given [1], [2], [6]. The central theme of this paper is to find the solutions of two integral equations with the G-function as kernel. As a large variety of functions that occur frequently in the problems of analysis and mathematical physics are only specialised or limiting forms of the kernel used in the present integral equations, our problem may prove of general interest. Later some specialised cases are derived from the main results.

We recall the definition of Meijer's G-function [4, p. 207]:

$$G = (2xi)^{-1} \begin{cases} x \begin{vmatrix} a_1, \dots, a_p \\ b_1, \dots, b_q \end{vmatrix} = G = G = M, n \\ p, q \\ x \begin{vmatrix} (a_p) \\ (b_q) \end{vmatrix} = G = G = G = G = G$$

$$= (2xi)^{-1} \begin{cases} \prod_{j=1}^{m} \Gamma(b_j + s) \prod_{j=1}^{n} \Gamma(1 - a_j - s) \\ \prod_{j=1}^{n} \Gamma(1 - b_j - s) \prod_{j=1}^{p} \Gamma(a_j + s) \end{cases}$$

$$= (2xi)^{-1} \begin{cases} \prod_{j=1}^{m} \Gamma(1 - b_j - s) \prod_{j=1}^{n} \Gamma(a_j + s) \\ \prod_{j=1}^{n} \Gamma(1 - b_j - s) \prod_{j=1}^{n} \Gamma(a_j + s) \end{cases}$$

$$= (2xi)^{-1} \begin{cases} \prod_{j=1}^{m} \Gamma(1 - b_j - s) \prod_{j=1}^{n} \Gamma(a_j + s) \\ \prod_{j=1}^{n} \Gamma(a_j + s) \end{bmatrix}$$

$$= (2xi)^{-1} \begin{cases} \prod_{j=1}^{m} \Gamma(a_j + s) \prod_{j=1}^{n} \Gamma(a_j + s) \\ \prod_{j=1}^{n} \Gamma(a_j + s) \prod_{j=1}^{n} \Gamma(a_j + s) \end{cases}$$

$$= (2xi)^{-1} \begin{cases} \prod_{j=1}^{m} \Gamma(a_j + s) \prod_{j=1}^{n} \Gamma(a_j + s) \\ \prod_{j=1}^{n} \Gamma(a_j + s) \end{bmatrix}$$

The following results will be needed in the development of the solutions.

A special case of [5, p. 212, (79)], taking  $\mu = 1$  and replacing  $\alpha$  by zy, is expressed as:

$$\int_{y}^{\infty} G \frac{3,1}{3,5} \left[ z(xy) \middle| \frac{k-m-1/2-\nu/2, -k+m+1/2+\nu/2, -1}{\nu/2-\lambda-m, \nu/2-\lambda+m, 0, -\nu/2+\lambda+m, -\nu/2+\lambda-m} \right] dx =$$

$$= y G \frac{2,1}{2,4} \left[ zy^{2} \middle| \frac{k-m-1/2-\nu/2, -k+m+1/2+\nu/2}{\nu/2-\lambda-m, \nu/2-\lambda+m, -\nu/2+\lambda+m, -\nu/2+\lambda-m} \right],$$

$$-R_{e}(k-m-1/2-\nu/2) > 0.$$
(1.2)

Similarly, another special case of [5, p. 200, (97)], taking  $\mu = 1$  and replacing  $\alpha$  by zy, is given by the relation:

$$\int_{0}^{y} G_{3,5}^{2,2} \left[ z(xy) \middle| \begin{array}{c} -1, k-m-1/2-\nu/2, -k+m+1/2+\nu/2 \\ \nu/2-\lambda-m, \nu/2-\lambda+m, -\nu/2+\lambda+m, -\nu/2+\lambda-m, 0 \end{array} \right] dx$$

$$= y G_{2,4}^{2,1} \left[ zy^{2} \middle| \begin{array}{c} k-m-1/2-\nu/2, -k+m+1/2+\nu/2 \\ \nu/2-\lambda-m, \nu/2-\lambda+m, -\nu/2+\lambda+m, -\nu/2+\lambda-m \end{array} \right],$$

$$R_{e}(\nu/2-\lambda-m) > -1, \quad R_{e}(\nu/2-\lambda+m) > -1. \quad (1.3)$$

# 2. Theorem 1.

If the integral equations

$$\int_{0}^{x} G \frac{3.1}{3.5} \left[ z(xy) \middle| \begin{array}{c} k - m - 1/2 - \nu/2, -k + m + 1/2 + \nu/2, -1 \\ \nu/2 - \lambda - m, \nu/2 - \lambda + m, 0, -\nu/2 + \lambda + m, -\nu/2 + \lambda - m \end{array} \right].$$

$$f(y) dy = \Phi(x, z),$$

$$-R_{e}(k - m - 1/2 - \nu/2) > 0,$$
(2.1)

and

$$\int_{0}^{\infty} \Phi(x, z) dx = \Psi(z)$$
 (2.2)

exist, then the solution of (2.1) is given as

$$f(y) = y^{-1} \int_{0}^{\infty} G \frac{2,1}{2,4} \left[ zy^{2} \middle| \begin{array}{c} k - m - 1/2 - \nu/2, -k + m + 1/2 + \nu/2 \\ \nu/2 - \lambda - m, \nu/2 - \lambda + m, -\nu/2 + \lambda + m, -\nu/2 + \lambda - m \end{array} \right].$$

$$. \Psi(z) dz. \qquad (2.3)$$

**Proof.** Let us suppose that both  $\Phi(x, z)$  and  $\Psi(z)$  exist. Substituting the value of  $\Phi(x, z)$  from (2.1) in (2.2), we obtain

$$\int_{0}^{\infty} \left\{ \int_{0}^{x} G \frac{3,1}{3,5} \left[ z(xy) \middle| \frac{k-m-1/2-\nu/2, -k+m+1/2+\nu/2, -1}{\nu/2-\lambda-m, \nu/2-\lambda+m, 0, -\nu/2+\lambda+m, -\nu/2+\lambda-m} \right] . f(y) dy \right\} dx = \Psi(z) .$$

By inverting the order of integration which permissible due to the convergence of the integrals involved, one can get

$$\int_{0}^{\infty} f(y) \left\{ G \frac{3.1}{3.5} \left[ z(xy) \middle| \frac{k - m - 1/2 - \nu/2, -k + m + 1/2 + \nu/2, -1}{\nu/2 - \lambda - m, \nu/2 - \lambda + m, 0, -\nu/2 + \lambda + m, -\nu/2 + \lambda - m} \right] \cdot x \, dx \right\} dy = \Psi(z) .$$

From (1.2), we arrive at

$$\int_{0}^{\infty} G \frac{2,1}{2,4} \left[ zy^{2} \middle| \begin{array}{c} k-m-1/2-\nu/2, -k+m+1/2+\nu/2 \\ \nu/2-\lambda-m, \nu/2-\lambda+m, -\nu/2+\lambda+m, -\nu/2+\lambda-m \end{array} \right] yf(y) dy =$$

$$= \Psi(z).$$

This is now in the form of the generalised Hankel transform [7, (1.1)], and hence inverting by applying [5, p. 5, (1)], we can find the solution (2.3).

# 3. Applications

(i) With  $\lambda = -m$ , our theorem leads to

Corollary 1. If the integral equations

$$\int_{0}^{x} G \frac{3.1}{3.5} \left[ [z(xy) \mid k-m-1/2-\nu/2, -k+m+1/2+\nu/2, -1] \right] f(y) dy = \frac{\nu/2, \nu/2+2m, 0, -\nu/2, -\nu/2-2m}{p_1(x, z), -R_e(k-m-1/2-\nu/2) > 0,}$$
(3.1)

and

$$\int_0^\infty \Phi_1(x,z) dx = \Psi_1(z) , \qquad (3.2)$$

exist, then the solution of (3.1) is given by

$$f(y) = y^{-1} \int_{0}^{\infty} G \frac{2,1}{2,4} \left[ zy^{2} \middle| \begin{array}{c} k - m - 1/2 - \nu/2, -k + m + 1/2 + \nu/2 \\ \nu/2, 2m + \nu/2, -\nu/2, -2m - \nu/2 \end{array} \right] \Psi_{1}(z) dz, (3.3)$$

where the generalised Hankel kernel in (3.3) is introduced by Bhise [3].

(ii) When  $\lambda = -m$ , k + m = 1/2 and using [4, p. 216, (3)] the theorem leads to Corollary 2. If the integral equations

$$\int_{0}^{x} G \frac{3.1}{3.5} \left[ z(xy) \middle| \frac{-2m-\nu/2, 2m+\nu/2, -1}{\nu/2, \nu/2+2m, 0, -\nu/2, -\nu/2-2m} \right] f(y) dy = \Phi_{2}(x,z), \\ -R_{e}(-2m-\nu/2) > 0,$$
 (3.4)

and

$$\int_0^\infty \Phi_2(x,z) dx = \Psi_2(z), \qquad (3.5)$$

exist, then the solution of (3.4) is expressed as

$$f(y) = y^{-1} \int_{0}^{\infty} \mathcal{F}(2yz^{1/2}) \Psi_2(z) dz$$
, (3.6)

where  $\mathcal{J}_v(x)$  is Bessel function.

# 4. Theorem 2.

If the integral equations

$$\int_{x}^{\infty} G \frac{2,2}{3,5} \left[ z(xy) \middle| \frac{-1, k-m-1/2-\nu/2, -k+m+1/2+\nu/2}{\nu/2-\lambda-m, \nu/2-\lambda+m, -\nu/2+\lambda+m, -\nu/2+\lambda-m, 0} \right].$$

$$f(y) dy = \Phi_{3}(x, z),$$

$$R_{e}(\nu/2-\lambda-m) > -1, \quad R_{e}(\nu/2-\lambda+m) > -1, \quad (4.1)$$

and

$$\int_0^\infty \Phi_3(x,z) dx = \Psi_3(z) , \qquad (4.2)$$

exist, then the solution of (4.1) is expressed as:

$$f(y) = y^{-1} \int_{0}^{\infty} G \frac{2,1}{2,4} \left[ zy^{2} \middle| \begin{array}{c} k - m - 1/2 - \nu/2, -k + m + 1/2 + \nu/2 \\ \nu/2 - \lambda - m, \nu/2 - \lambda + m, -\nu/2 + \lambda + m, -\nu/2 + \lambda - m \end{array} \right].$$

$$\cdot \Psi_{3}(z) dz. \qquad (4.3)$$

The proof of this theorem easily follows by proceeding on the lines of Theorem. 1.

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