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Katedra algebry a geometrie přírodovědecké fakulty Univerzity Palackého v Olomouci Vedoucí katedry: prof. RNDr. L. Sedláček, CSc.

TWO APPLICATIONS OF AN INTEGRAL FORMULA

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We are going to present two consequences of a general integral formula presented in [1].

1. Harmonic mappings of Riemannian manifolds

Be given a Riemannian manifold (M, ds^2) , dim M = m. In a suitable domain $U \subset M$, let us write (i, j, ... = 1, ..., m)

$$ds^2 = \sum_i (\omega^i)^2, \tag{1.1}$$

 ω^i being linearly independent 1-forms on U. Then there are on U 1-forms ω_i^j such that

$$d\omega^{i} = \sum_{i} \omega^{j} \wedge \omega_{j}^{i}, \qquad \omega_{i}^{j} + \omega_{i}^{j} = 0;$$
(1.2)

the forms ω_i^j are uniquely determined by (1.2). The components of the curvature tensor be introduced by

$$d\omega_i^j = \sum_k \omega_i^k \wedge \omega_k^j - \frac{1}{2} R_{ikl}^j \omega^k \wedge \omega^l, \qquad R_{ikl}^j + R_{ilk}^j = 0; \tag{1.3}$$

they satisfy the symmetry relations

$$R_{ikl}^{j} + R_{ikl}^{i} = 0, \qquad R_{ikl}^{j} = R_{kij}^{l}, \qquad R_{ikl}^{j} + R_{ilj}^{k} + R_{ijk}^{l} = 0.$$
 (1.4)

Let $v_1, ..., v_m$ be the field of orthonormal frames on U dual to the field of coframes $\omega^1, ..., \omega^m$. Denote by $K(v_i, v_j)$, $i \neq j$, the sectional curvature of the 2-plane $\{v_i, v_j\}$; of course, $K(v_i, v_j) = R_{iij}^j$.

Further, be given another Riemannian manifold $(N, d\sigma^2)$, dim N = n, and a mapping $f: M \to N$. Consider a neighbourhood $V \subset N$ such that $f(U) \subset V$ and there are 1-forms φ^a $(\alpha, \beta, \ldots = 1, \ldots, n)$ satisfying

$$d\sigma^2 = \sum (\varphi^{\alpha})^2. \tag{1.5}$$

Let

$$\tau^{\alpha} := f^* \varphi^{\alpha} = \sum_{i} A^{\alpha}_{i} \omega^{i}, \qquad \tau^{\beta}_{\alpha} := f^* \varphi^{\beta}_{\alpha}. \tag{1.6}$$

The exterior differentiation of (1.6₁) yields

$$\sum_{i} (\mathrm{d}A_{i}^{\alpha} - \sum_{j} A_{j}^{\alpha} \omega_{i}^{j} + \sum_{\beta} A_{i}^{\beta} \tau_{\beta}^{\alpha}) \wedge \omega^{i} = 0, \tag{1.7}$$

and, according to E. Cartan's lemma, we get the existence of functions A_{ij}^{α} on U satisfying

$$dA_i^{\alpha} - \sum_j A_j^{\alpha} \omega_i^j + \sum_{\beta} A_i^{\beta} \tau_{\beta}^{\alpha} = \sum_j A_{ij}^{\alpha} \omega^j, \qquad A_{ij}^{\alpha} = A_{ji}^{\alpha}.$$
 (1.8)

A further exterior differentiation implies

$$\sum_{j} (\mathrm{d}A_{ij}^{\alpha} - \sum_{k} A_{ik}^{\alpha} \omega_{j}^{k} - \sum_{k} A_{kj}^{\alpha} \omega_{i}^{k} + \sum_{\beta} A_{ij}^{\beta} \tau_{\beta}^{\alpha}) \wedge \omega^{j} =$$

$$= \frac{1}{2} \sum_{j,k} (\sum_{l} A_{l}^{\alpha} R_{ijk}^{l} - \sum_{\beta,\gamma,\delta} A_{i}^{\beta} A_{j}^{\gamma} A_{k}^{\delta} S_{\beta\gamma\delta}^{\alpha}) \omega^{j} \wedge \omega^{k}, \qquad (1.9)$$

 $S_{\beta\gamma\delta}^{\alpha}$ being the components of the curvature tensor of $(N, d\sigma^2)$. Thus there are functions A_{ijk}^{α} such that

$$dA_{ij}^{\alpha} - \sum_{k} A_{ik}^{\alpha} \omega_{j}^{k} - \sum_{k} A_{kj}^{\alpha} \omega_{i}^{k} + \sum_{\beta} A_{ij}^{\beta} \tau_{\beta}^{\alpha} = \sum_{k} A_{ijk}^{\alpha} \omega^{k}, \qquad A_{ijk}^{\alpha} = A_{jik}^{\alpha}, \quad (1.10)$$

$$A_{ijk}^{\alpha} - A_{ikj}^{\alpha} = \sum_{l} A_{l}^{\alpha} R_{ikj}^{l} - \sum_{\beta,\gamma,\delta} A_{\delta}^{\beta} A_{k}^{\gamma} A_{\delta}^{\delta} S_{\beta\gamma\delta}^{\alpha}. \tag{1.11}$$

Let us consider on U the 1-forms

$$\varphi_1 = \sum_{\alpha, i, j} A_i^{\alpha} A_{ij}^{\alpha} \omega^j, \qquad \varphi_2 = \sum_{\alpha, i, j} A_j^{\alpha} A_{ii}^{\alpha} \omega^j.$$
(1.12)

It is easy to see that the forms (1.12) are globally defined over all of M. The usual \star -operator be defined by

$$\star \omega^{i} = (-1)^{i+1} \omega^{1} \wedge \dots \wedge \omega^{i-1} \wedge \omega^{i+1} \wedge \dots \wedge \omega^{n},$$
i.e., $d\sigma := \omega^{1} \wedge \dots \wedge \omega^{n} = \omega^{i} \wedge \star \omega^{i}.$

$$(1.13)$$

Now,

$$d * \varphi_1 = \sum_{\alpha, i, j} \{ (A_{ij}^{\alpha})^2 + A_i^{\alpha} A_{jij}^{\alpha} \} do,$$

$$d * \varphi_2 = \sum_{\alpha, i, j} (A_{ii}^{\alpha} A_{jj}^{\alpha} + A_i^{\alpha} A_{jji}^{\alpha}) do,$$
(1.14)

and, according to (1.11),

$$d * (\varphi_1 - \varphi_2) = \sum_{\alpha, i, j} \{ (A_{ij}^{\alpha})^2 - A_{ii}^{\alpha} A_{jj}^{\alpha} + \sum_k A_i^{\alpha} A_k^{\alpha} R_{jji}^k - \sum_{\beta, \gamma, \delta} A_i^{\alpha} A_i^{\delta} A_i^{\gamma} A_j^{\delta} S_{\beta\gamma\delta}^{\alpha} \} do. \quad (1.15)$$

Let us turn our attention to the geometrical interpretation of the above introduced invariants. Let $p \in U \subset M$ be a given point. The Euclidean connection on M or N resp. is given by

$$\nabla m = \sum_{i} \omega^{i} v_{i}, \qquad \nabla v_{i} = \sum_{j} \omega_{i}^{j} v_{j} \qquad \text{or}$$

$$\nabla^{*} n = \sum_{\alpha} \varphi^{\alpha} w_{\alpha}, \qquad \nabla^{*} w_{\alpha} = \sum_{\beta} \varphi^{\beta}_{\alpha} w_{\beta} \qquad \text{resp.};$$
(1.16)

here, w_1, \ldots, w_n is the dual basis to $\varphi^1, \ldots, \varphi^n$. Evidently,

$$\mathrm{d}f_{\mathbf{r}}(v_i) = A_i^{\alpha} w_{\alpha}. \tag{1.17}$$

Let $v \in T_p(M)$ be a non-zero vector. Choose a curve $\gamma: (-\varepsilon, \varepsilon) \to M$ such that $\gamma(0) = p$; let s be its arc and v its tangent vector at p. Denote by $\gamma^* = f \circ \gamma: (-\varepsilon, \varepsilon) \to N$ the corresponding curve. Then it is easy to see that

$$\frac{\nabla^* n}{\mathrm{d}s^2} - \mathrm{d}f_p \left(\frac{\nabla^2 m}{\mathrm{d}s^2} \right) = \frac{L(v)}{|v|^2} \,, \tag{1.18}$$

where $|v|^2 = \sum_{i} (\omega^{i}(v))^2$ and

$$L(v) = A_{ii}^{\alpha}(p) \,\omega^{i}(v) \,\omega^{j}(v) \,w_{\sigma}(f(p)). \tag{1.19}$$

This gives the geometrical interpretation of the quadratic mapping

$$L: T_p(M) \to T_{f(p)}(N).$$
 (1.20)

Let L(.,.) be the corresponding bilinear mapping.

At p, let us choose an orthonormal frame v_i , let w_x be an orthonormal frame at f(p). Then

$$L(v_i, v_j) = A_{ij}^{\alpha} w_{\alpha} \tag{1.21}$$

and the expressions

$$\sum_{i,j} |L(v_i, v_j)|^2 = \sum_{i,j,\alpha} (A_{ij}^{\alpha})^2, \qquad |\sum_{i} L(v_i)|^2 = \sum_{\alpha,i,ji} A_{ii}^{\alpha} A_{jj}^{\alpha}$$
(1.22)

do not depend on the choice of the frames v_i and w_a . In the same way, the vector

$$t = \sum_{i} L(v_i) \tag{1.23}$$

is invariant; the mapping

$$t: M \to T(N), \qquad t(p) \in T_{f(p)}(N)$$
 (1.24)

is the so-called tension field. The mapping $f: M \to N$ is said to be harmonic if t = 0 for each $p \in M$.

The frames (v_1, \ldots, v_m) and (w_1, \ldots, w_n) at p and f(p) resp. are called adapted to f if

$$\begin{aligned} \mathrm{d}f_p(v_i) &= A_i w_i & \text{for } i = 1, \dots, m \text{ in the case } m \leq n \text{ and} \\ \mathrm{d}f_p(v_\alpha) &= A_\alpha w_\alpha & \text{for } \alpha = 1, \dots, n, \\ \mathrm{d}f_p(v_\varrho) &= 0 & \text{for } \varrho = n + 1, \dots, m \text{ in the case } m > n. \end{aligned} \tag{1.25}$$

Thus, we may always write (25_1) setting $w_i = 0$ for i > n. The adapted bases exist for each couple (p, f(p)). In the adapted bases, we have

$$\sum_{\alpha,i,j} \sum_{k} A_{i}^{\alpha} A_{k}^{\alpha} R_{jji}^{k} = \sum_{i} (A_{i})^{2} \sum_{j \neq i} K(v_{j}, v_{i}),$$
 (1.26)

$$\sum_{\alpha,l,j} \sum_{\beta,\gamma,\delta} A_j^{\alpha} A_i^{\beta} A_i^{\gamma} A_j^{\delta} S_{\beta\gamma\delta}^{\alpha} = 2 \sum_{i \neq j} (A_i A_j)^2 K^*(w_i, w_j). \tag{1.27}$$

Further.

$$\varphi_1(v_i) = \sum_i \langle \mathrm{d}f(v_j), L(v_i, v_j) \rangle, \qquad \varphi_2(v_i) = \langle \mathrm{d}f(v_i), t \rangle,$$
 (1.28)

 \langle , \rangle being the scalar product in $T_{f(p)}(N)$.

Choosing for each couple (p, f(p)) the adapted bases, we have the integral formula

$$\int_{\partial M} * (\varphi_1 - \varphi_2) = \int_{M} \left\{ \sum_{i,j} |L(v_i, v_j)|^2 - |t|^2 + \sum_{i} (A_i)^2 \sum_{j \neq i} K(v_j, v_i) - 2 \sum_{i \neq j} (A_i A_j)^2 K^*(w_i, w_j) \right\} dv.$$
 (1.29)

Thus we get the following

Theorem. Let M, N be Riemannian manifolds and $f: M \to N$ a harmonic mapping. Let N have non-positive sectional curvatures and let M have, at each point $p \in M$ and for each unit vector $v \in T_p(M)$ the following property: v_1, \ldots, v_{m-1}, v being an orthonormal basis of $T_p(M)$, we have $\sum_{r=1,\ldots,m-1} K(v,v_r) > 0$. Let $\varphi_1 = \varphi_2$ on the boundary ∂M of M. Then f is a constant mapping.

2. Holomorphic curves in the Hermitian plane

Be given a Hermitian plane H^2 and let $m: D \to H^2$ be a holomorphic curve, $D \subset \mathscr{C}$ being a bounded domain. To each its point m(d), $d \in D$, let us associate an orthonormal frame $\{m, w_1, w_2\}$. Then we have the equations

$$dm = \tau^1 w_1 + \tau^2 w_2,$$

$$dw_1 = \tau_1^1 w_1 + \tau_1^2 w_2, \qquad dw_2 = \tau_2^1 w_1 + \tau_2^2 w_2;$$
(2.1)

clearly (i, j, ... = 1, 2)

$$\tau_i^j + \bar{\tau}_i^i = 0, \tag{2.2}$$

$$d\tau^{i} = \tau^{j} \wedge \tau^{i}_{j}, \qquad d\tau^{j}_{i} = \tau^{k}_{i} \wedge \tau^{j}_{k}. \tag{2.3}$$

Let us restrict ourselves to the tangent frames satisfying

$$\tau^2 = 0. \tag{2.4}$$

By successive exterior differentiations we get the existence of functions $R, S, T, U: D \to \mathcal{C}$ such that

$$\tau_1^2 = R\tau^1, \tag{2.5}$$

$$dR + R(\tau_2^2 - 2\tau_1^1) = S\tau^1, \tag{2.6}$$

$$dS + S(\tau_2^2 - 3\tau_1^1) + 3R^2 \bar{R}_{\tau}^{-1} = T\tau^1, \tag{2.7}$$

$$dT + T(\tau_2^2 - 4\tau_1^1) + 10R\bar{R}S\bar{\tau}^1 = U\tau^1.$$
 (2.8)

Let us consider another field of orthonormal frames

$$u_1 = e^{i\alpha}w_1, \qquad u_2 = e^{i\beta}w_2; \qquad \alpha, \beta: D \to \mathcal{R};$$
 (2.9)

let us write

$$dm = \varphi^1 u_1, \quad du_1 = \varphi_1^1 u_1 + \varphi_1^2 u_2, \quad du_2 = \varphi_2^1 u_1 + \varphi_2^2 u_2.$$
 (2.10)

Then it is easy to see that

$$\varphi^1 = e^{-i\alpha}\tau^1,\tag{2.11}$$

$$\varphi_1^1 = \tau_1^1 + i \, d\alpha, \qquad \varphi_2^2 = \tau_2^2 + i \, d\beta, \qquad \varphi_1^2 = e^{i(\alpha - \beta)} \tau_1^2.$$
 (2.12)

Write

$$\varphi_1^2 = R'\varphi_1, \tag{2.13}$$

$$dR' + R'(\varphi_2^2 - 2\varphi_1^1) = S'\varphi^1, \tag{2.14}$$

$$dS' + S'(\varphi_2^2 - 3\varphi_1^1) + 3R'^2 \bar{R}' \bar{\varphi}^1 = T' \varphi^1,$$
(2.15)

$$dT' + T'(\varphi_2^2 - 4\varphi_1^1) + 10R'\bar{R}'S'\bar{\varphi}^1 = U'\varphi^1.$$
 (2.16)

Then

$$R' = e^{i(2\alpha - \beta)}R,$$
 $S' = e^{i(3\alpha - \beta)}S,$
 $T' = e^{i(4\alpha - \beta)}T.$ $U' = e^{i(5\alpha - \beta)}U.$ (2.17)

The mappings $B^{(k)}: T_m \to N_m$ be introduced by

$$B(zw_1) = z^2 Rw_2, B^{(1)}(zw_1) = z^3 Sw_2, B^{(2)}(zw_1) = z^4 Tw_2, z \in \mathscr{C}. (2.18)$$

These mappings are invariant. Indeed: Let $w = zw_1 = z'u_1$, then $z' = e^{-i\alpha}z$ and $z'^2R'u_2 = z^2Rw_2$; similarly for $B^{(k)}$. Let $S^1 = \{w \in T_m; \langle w, w \rangle = 1\}$, i.e., $S^1 = \{zw_1; |z|^2 = 1\}$. Then $B^{(k)}(S^1)$ is a circle; the radius of $B^{(1)}(S^1)$ is equal to $|S|^{1/2}$, the radius of $B^{(1)}(S^1)$ is equal to $|S|^{1/2}$, etc. The geometrical interpretation of the mappings $B^{(k)}$ will be presented later on.

The area element of m is given by

$$do = \frac{1}{2}i\tau^1 \wedge \overline{\tau}^1. \tag{2.19}$$

The Hodge operator be introduced by

$$\times \tau^1 = -i\tau^1, \qquad \times \bar{\tau}^1 = i\bar{\tau}^1. \tag{2.20}$$

Let $f: D \to \mathcal{R}$ be a function. Then its Laplacian Δf is given, as usually, by

$$\Delta f \, \mathrm{d}o = \mathrm{d} \star \mathrm{d}f. \tag{2.21}$$

The straightforward calculations lead to $(n \ge 1)$

$$d | R |^{2n} = 2n | R |^{2n-2} Re (\bar{R}S\tau^1),$$
 (2.22)

$$\Delta \mid R \mid^{2n} = 4n \mid R \mid^{2n-2} (n \mid S \mid^2 - 3 \mid R \mid^4), \tag{2.23}$$

$$d |S|^{2n} = 2n |S|^{2n-2} Re \{ (\bar{S}T - 3SR\bar{R}^2) \tau^1 \}, \qquad (2.24)$$

$$\Delta |S|^{2n} = 4n\{ |S|^{2n-2}(n|T|^2 - 16|S|^2|R|^2 + 9n|R|^6) - 6(n-1)|S|^{2n-4}|R|^2 \operatorname{Re}(\bar{S}^2RT) \}.$$
(2.25)

Especially,

$$\Delta |R|^{2} = 4(|S|^{2} - 3|R|^{4}),$$

$$\Delta |R|^{4} = 8|R|^{2}(2|S|^{2} - 3|R|^{4}),$$

$$\Delta |S|^{2} = 4(|T|^{2} - 16|S|^{2}|R|^{2} + 9|R|^{6})$$
(2.26)

and

$$\Delta(|S|^2 + 4|R|^4) = 4(|T|^2 - 15|R|^6). \tag{2.27}$$

Lemma. Let S = 0 on D. Then m(D) is a part of a straight line of H^2 . Proof. The equation (2.7) implies R = 0. QED.

Theorem. Let S = 0 on ∂D and

$$3 |R|^4 \ge |S|^2$$
 on D . (2.28)

Then m(D) is a part of a straight line of H^2 .

Proof. Obviously, $\star d \mid R \mid^{2n} = 0$ on ∂D . From the integral formula

$$0 = \int_{M} \Delta |R|^2 dv \qquad (2.29)$$

we get, because of (2.26),

$$3 |R|^4 = |S|^2$$
 on D . (2.30)

The integral formula

$$\int_{\partial M} d |R|^4 = 8 \int_{M} |R|^2 (2 |S|^2 - 3 |R|^4) dv$$
 (2.31)

reduces to

$$0 = \int_{M} |R|^{6} dv, \qquad (2.32)$$

and we get R = 0. OED.

The formulas (2.23), (2.25), (2.27) imply new characterizations of straight lines of H^2 . It is sufficient to suppose S = 0 on ∂D and, for ex.,

$$14 | R |^6 \ge | T |^2 \tag{2.33}$$

or

$$|T|^2 \ge 8 |R|^2 (2|S|^2 - |R|^4)$$
 (2.34)

on D; see (2.27) and (2.26₃).

Now, the geometrical description of the mappings $B^{(k)}$ is given in [1]. To do this, let us consider H^2 as a space over m, i.e., H^2 becomes E^4 . Write

$$v_{1} = w_{1}, v_{2} = iw_{1}, v_{3} = w_{2}, v_{4} = iw_{2},$$

$$\tau^{1} = \omega^{1} + i\omega^{2}, \tau^{2} = \omega^{3} + i\omega^{4}, \tau_{1}^{2} = \omega_{1}^{3} + i\omega_{1}^{4},$$

$$\tau_{1}^{1} = i\omega_{1}^{2}, \tau_{2}^{2} = i\omega_{3}^{4},$$

$$(2.35)$$

i.e.,

$$dm = \omega^{1}v_{1} + \omega^{2}_{2},$$

$$dv_{1} = \omega^{2}_{1}v_{2} + \omega^{3}_{1}v_{3} + \omega^{4}_{1-4},$$

$$dv_{2} = -\omega^{2}_{1}v_{1} - \omega^{4}_{1}v_{3} + \omega^{3}_{1}\iota_{4},$$

$$dv_{3} = -\omega^{3}_{1}v_{1} + \omega^{4}_{1}v_{2} + \omega^{4}_{3}v_{4},$$

$$dv_{4} = -\omega^{4}_{1}v_{1} - \omega^{3}_{1}v_{2} - \omega^{4}_{3}v_{3}$$

$$(2.36)$$

and

$$\omega_1^3 = R_1 \omega^1 - R_2 \omega^2, \qquad \omega_1^4 = R^2 \omega^1 + R_1 \omega^2$$
 (2.37)

with $R_1 = \text{Re } R$, $R_2 = \text{Im } R$. In E^4 , consider a general surface

$$dn = \varrho^{1}v_{1} + \varrho^{2}v_{2},$$

$$dv_{1} = \varrho_{1}^{2}v_{2} + \varrho_{1}^{3}v_{3} + \varrho_{1}^{4}v_{4},$$

$$dv_{2} = -\varrho_{1}^{2}v_{1} + \varrho_{2}^{3}v_{3} + \varrho_{2}^{4}v_{4},$$

$$dv_{3} = -\varrho_{1}^{3}v_{1} - \varrho_{2}^{3}v_{2} + \varrho_{3}^{4}v_{4},$$

$$dv_{4} = -\varrho_{1}^{4}v_{1} - \varrho_{2}^{4}v_{2} - \varrho_{3}^{4}v_{3}$$

$$(2.38)$$

with

$$\varrho_{1}^{3} = a_{1}\varrho^{1} + a_{2}\varrho^{2}, \qquad \varrho_{2}^{3} = a_{2}\varrho^{1} + a_{3}\varrho^{2},
\varrho_{1}^{4} = b_{1}\varrho^{1} + b_{2}\varrho^{2}, \qquad \varrho_{2}^{4} = b_{2}\varrho^{1} + b_{3}\varrho^{2},
da_{1} - 2a_{2}\varrho_{1}^{2} - b_{1}\varrho_{3}^{4} = \alpha_{1}\varrho^{1} + \alpha_{2}\varrho^{2}.$$
(2.39)

$$da_{2} + (a_{1} - a_{3}) \varrho_{1}^{2} - b_{2}\varrho_{3}^{4} = \alpha_{2}\varrho^{1} + \alpha_{3}\varrho^{2},$$

$$da_{3} + 2a_{2}\varrho_{1}^{2} - b_{3}\varrho_{3}^{4} = \alpha_{3}\varrho^{1} + \alpha_{4}\varrho_{2},$$

$$db_{1} - 2b_{2}\varrho_{1}^{2} + a_{1}\varrho_{3}^{4} = \beta_{1}\varrho^{1} + \beta_{2}\varrho^{2},$$

$$db_{2} + (b_{1} - b_{3}) \varrho_{1}^{2} + a_{2}\varrho_{3}^{4} = \beta_{2}\varrho^{1} + \beta_{3}\varrho^{2},$$

$$db_{3} + 2b_{2}\varrho_{1}^{2} + a_{3}\varrho_{3}^{4} = \beta_{3}\varrho^{1} + \beta_{4}\varrho^{2}.$$

$$(2.40)$$

Then it is known [1] that, for

$$\Phi = (a_1\alpha_3 + a_2\alpha_4 - a_2\alpha_2 - a_3\alpha_3 + b_1\beta_3 + b_2\beta_4 - b_2\beta_2 - b_3\beta_3) \varrho^1 + (a_2\alpha_1 + a_3\alpha_2 - a_1\alpha_2 - a_2\alpha_3 + b_2\beta_1 + b_3\beta_2 - b_1\beta_2 - b_2\beta_3) \varrho^2,$$

we have

$$\int_{\partial N} \Phi = \int_{N} \left[2(\alpha_{1}\alpha_{3} + \alpha_{2}\alpha_{4} - \alpha_{2}^{2} - \alpha_{3}^{2} + \beta_{1}\beta_{3} + \beta_{2}\beta_{4} - \beta_{2}^{2} - \beta_{3}^{2}) - \right. \\ \left. - \left. \left\{ (a_{1} - a_{3})^{2} + 4a_{2}^{2} + (b_{1} - b_{3})^{2} + 4b_{2}^{2} \right\} (a_{1}a_{3} - a_{2}^{2} + b_{1}b_{3} - b_{2}^{2}) + \right. \\ \left. + 2\left\{ b_{2}(a_{1} - a_{3}) + a_{2}(b_{3} - b_{1}) \right\}^{2} \right] dv.$$
(2.42)

In our case, (2.42) is identical with

$$\int_{\partial N} * d |R|^2 = \int_{N} \Delta |R|^2 dv.$$
 (2.43)

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SOUHRN

DVĚ APLIKACE JEDNÉ INTEGRÁLNÍ FORMULE

ALOIS ŠVEC

V práci jsou vyloženy aplikace integrální formule [1] na teorii harmonických zobrazení a na teorii křivky v hermiteovské rovině.

РЕЗЮМЕ

ДВА ПРИМЕНЕНИЯ ОДНОЙ ИНТЕГРАЛЬНОЙ ФОРМУЛЫ

АЛОИС ШВЕЦ

В работе излагаются приложения интегральной формулы [1] на теорию гармонических отображений и на теорию кривых в пространствах Эрмита.