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PARALLEL DISPLACEMENT OF VECTORS ON RHEONOMOUS ANHOLONOMIC MANIFOLD

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We shall make use of the notations of [2]. Let us suppose that in the stationary space $r - L_n(t)$ there is defined the rheonomous anholonomic manifold $r - L_n^m(t)$ by means of the equations

(1)
$$B_a^{\alpha} = B_a^{\alpha}(x^{\omega}, t), \quad {}_{p}n^{\alpha} = {}_{p}n^{\alpha}(x^{\omega}, t), \quad x^{\alpha} = x^{\alpha}(u^{A}, t)$$

where $\{x^{\omega}, t\} \in \Omega$, $\{u^{A}, t\} \in \Lambda \times I$ (see [2], Equations (1,2) and (1,3)). Let the parametric equations

(2)
$$x^{\alpha} = x^{\alpha}(T), \quad t = T, \quad T \in J$$

describe a trajectory which lies on the rheonomous manifold $r - L_n^m(t)$. Thus for every $T \in J$ there exist m numbers du^a/dT such that the equation

(3)
$$\frac{\mathrm{d}x^{\alpha}}{\mathrm{d}T} = \frac{\mathrm{d}u^{a}}{\mathrm{d}T}B_{a}^{\alpha} + B_{t}^{\alpha}$$

holds.

Let $v^a(t)$, $w^a(t)$ be the functions of the class C_2 defined on the interval J. We shall suppose that along the trajectory (2) a field of virtual or tangential vectors of the rheonomous manifold $r - L_n^m(t)$ is defined by means of functions

$$(4) v^{\alpha} = v^{a}B_{a}^{\alpha}, \quad T \in J$$

or

(5)
$$w^{\alpha} = w^{a}B_{a}^{\alpha} + B_{t}^{\alpha}, \quad T \in J,$$

respectively. The vector field (4) or (5) will be called the W-parallel or T-parallel if for every $T \in J$

$$B^a_{\alpha}D_Tv^{\alpha}=0$$

or

$$B^a_{\alpha}D_Tw^{\alpha}=0,$$

respectively. By means of formulae (2,1) and (2,2) of [2] we shall easily calculate that

$$D_T(v^a B_a^x) = B_a^x (D_T v^a + w_b^a v^b) + {}_p n^x \left(h_{ab}^p v^b \frac{\mathrm{d} u^a}{\mathrm{d} T} + {}^p m_a v^a \right)$$

or

$$D_{T}(w^{a}B_{a}^{x} + B_{t}^{x}) = B_{a}^{x} \left(D_{T}w^{a} + w_{b}^{a}w^{b} + 'w_{b}^{a} \frac{\mathrm{d}u^{b}}{\mathrm{d}T} + W^{a} \right) +$$

$$+ {}_{p}n^{x} \left(h_{ab}^{p}w^{b} \frac{\mathrm{d}u^{a}}{\mathrm{d}T} + {}^{p}m_{a}w^{a} + {}^{p'}m_{a} \frac{\mathrm{d}u^{a}}{\mathrm{d}T} + {}^{p}W \right).$$

Hence and from (4) and (5) it follows that the equations (6) or (7) are equivalent to the equations

$$(8) D_T v^a + w_b^a v^b = 0$$

or

(9)
$$D_T w^a + w_b^a w^b + 'w_b^a \frac{\mathrm{d}u^b}{\mathrm{d}T} + W^a = 0,$$

respectively.

If for every $T \in J$

$$(10) D_T v^a = 0$$

then the vector field (4) will be called the pseudoparallel field.

We have defined three "parallel" displacements, let us first consider the *T*-parallel displacement. We shall prove the following theorem:

Let $[x^a(_{\circ}T),_{\circ}T]$ be a point of the trajectory (2) and $(_{\circ}w^a_{\circ}B^a_a + _{\circ}B^a_t)$ the tangential vector of the rheonomous manifold $r - L_n^m(t)$ defined at that point. Then there exists along the trajectory (2) exactly one T-parallel field (5) such that $w^a(_{\circ}T) = _{\circ}w^a$.

The proof is easy. The functions $w^a(T)$ are solutions of the system of *m*-differential equations of (9). Writing in full the system (9) we obtain the equivalent linear system of *m*-differential equations of *m*-unknown functions in Cauchy's canonical form. From this result it follows immediately that the system (9) possesses on the interval J exactly one solution for given initial conditions.

If the tangential vectors of a given trajectory of the rheonomous manifold $r - L_n^m(t)$ form the *T*-parallel field then we call such trajectory a *T*-geodesic. The following theorem holds:

Let $({}_{\circ}w^{\alpha}) = ({}_{\circ}w^{\alpha}{}_{\circ}B^{\alpha}_{a} + {}_{\circ}B^{\alpha}_{t})$ be the tangential vector of the rheonomous manifold defined at its point $[{}_{\circ}x^{\alpha}, {}_{\circ}T]$. Then there exists (locally) exactly one T-geodesic with

parametric description (2) such that

(11)
$$x^{\alpha}(_{\circ}T) = _{\circ}x^{\alpha}, \quad \frac{\mathrm{d}x^{\alpha}(_{\circ}T)}{\mathrm{d}T} = _{\circ}v^{\alpha} \equiv _{\circ}v^{\alpha}{_{\circ}}B^{\alpha}_{a} + _{\circ}B^{\alpha}_{t}.$$

Proof. By (9) the W-geodesic is described by the system of differential equations

(12)
$$D_T \frac{du^a}{dT} + (w_b^a + 'w_b^a) \frac{du^a}{dT} + W^a = 0, \quad \frac{dx^a}{dT} = B_a^a \frac{du^a}{dT} + B_t^x.$$

If we introduce the notation $\xi^a = du^a/dT$ then (12) may be easely written in Cauchy's canonical form. For initial conditions (11), now of the form $x^{\alpha}(_{\circ}T) = _{\circ}x^{\alpha}$, $\xi^a(_{\circ}T) = _{\circ}w^a$, the new equivalent system has (locally) exactly one solution.

Preceding consideration of the T-parallel displacement may be extended also on the cases of the W-parallel and pseudoparallel displacement. We may show that under "usual" conditions the virtual vector may undergo a W-parallel or pseudoparallel displacement along the given trajectory in exactly one way. We may also define the notion of the W-geodesic or the pseudogeodesic and show its unique (local) existence for usual initial conditions. In case when $r - L_n^m(t)$ is a stationary manifold then $w_h^a = w_h^a = 0$, $w_h^a = 0$, $w_h^a = 0$ and all three "parallel" displacement are identical.

If a metric field $(g_{\alpha\beta})$ is defined in L_n (everywhere symmetric and positively definite) then L_n is a *Riemannian* space. We shall denote it R_n . In this case we shall denote $r-L_n(t)$ by $r-R_n(t)$. In the space $r-R_n(t)$ let be given such rheonomous manifold $r-L_n^m(t)$ that for every $\{x^{\omega}, t\}$ $g_{\alpha\beta}B_{\alpha p}^{\alpha}n^{\beta}=0$, $g_{\alpha\beta}p^{\alpha}n^{\alpha}p^{\beta}=p_{\alpha}\delta$. Then we shall denote $r-L_n^m(t)$ by $r-R_n^m(t)$. By means of the equations

(13)
$$g_{ab} = g_{\alpha\beta}B_a^{\alpha}B_b^{\beta}, \quad g_{ab}g^{bc} = \delta_a^c$$

the coordinates of the metric tensor field (g_{ab}) are defined on the rheonomous manifold $r - R_n^m(t)$.

We shall show that the following equations hold at every point of the rheonomous manifold $r - R_n^m(t)$:

$$(14) D_c g_{ba} = 0,$$

(15)
$$\Gamma_{cb}^{\ a} = \begin{Bmatrix} a \\ cb \end{Bmatrix} + \Omega_{bc}^{\ a} + \Omega_{cb}^{\ a} + \Omega_{bc}^{\ a}$$

where

$$\begin{cases} a \\ cb \end{cases} = \frac{1}{2} g^{ae} \big(\partial_c g_{be} \, + \, \partial_b g_{ce} \, - \, \partial_e g_{cb} \big) \, , \quad \Omega_{bc}^{\ a} = \, \partial_{[b} B^\alpha_{|a|} B^\alpha_{c]} \, , \label{eq:cb}$$

$$D_t g_{\beta\alpha} = 0.$$

I. From the equation

(17)
$$g_{\alpha\beta} = g_{\gamma\delta}\delta^{\gamma}_{\alpha}\delta^{\delta}_{\beta} = g_{ab}B^{a}_{\alpha}B^{b}_{\beta} + {}_{pq}\delta^{p}n_{\alpha}{}^{q}n_{\beta}$$

and the relation $\nabla_{\gamma}g_{\beta\alpha}=0$ it follows that

$$D_c g_{ba} = B_c^{\gamma} B_b^{\beta} B_a^{\alpha} \nabla_{\gamma} (B_b^{\epsilon} B_a^{f} g_{\epsilon f}) = B_c^{\gamma} B_b^{\beta} B_a^{\alpha} \nabla_{\gamma} (g_{\beta \alpha} - {}_{na} \delta^{p} n_{\beta}^{q} n_{\alpha}) = 0.$$

II. Obviously we may write

$$\Gamma_{cb}{}^{a} = \frac{1}{2} B_{c}^{y} B_{b}^{\beta} B_{\alpha}^{a} \left[\partial_{y} \left(B_{\beta}^{e} B_{\omega}^{g} g_{eg} + {}_{pq} \delta^{p} n_{\beta}^{q} n_{\omega} \right) + \right. \\ \left. + \left. \partial_{\beta} \left(B_{\gamma}^{f} B_{\omega}^{g} g_{fg} + {}_{pq} \delta^{p} n_{\gamma}^{q} n_{\omega} \right) + \partial_{\omega} \left(B_{\gamma}^{f} B_{\beta}^{e} g_{fe} + {}_{pq} \delta^{p} n_{\gamma}^{q} n_{\beta} \right) \right] g^{\alpha \omega} - B_{b}^{\beta} \partial_{c} B_{\beta}^{a} .$$

Applying a simple modification by means of (17) we obtain (15).

III. Let us denote by the symbol ∇_t the absolute derivate in the space R_n along the curve which is described by the parametric equations $x^{\alpha} = x^{\alpha}(_{\circ}u^{A}, t)$, $t \in I$. Evidently $D_t g_{\alpha\beta} = \nabla_t g_{\alpha\beta} = 0$.

We shall show that the following theorem holds: Let (u^a) , (v^b) be two fields of W-parallel vectors which are defined along the trajectory (2). Then

(18)
$$D_T(g_{ab}u^av^b) = 0^{1}$$

at every point of the trajectory (2).

In the proof of the theorem we shall make use of the notation $G_{ab} = D_t g_{ab}$. From (13) and (2,4) in [2] it follows that

$$G_{ab} = g_{\alpha\beta}(D_t B_a^{\alpha}) B_b^{\beta} + g_{\alpha\beta} B_a^{\alpha} D_t B_b^{\beta} = w_{ab} + w_{ba}.$$

Hence and from the equations

$$D_T u^a \, + \, w^a_b u^b \, = \, 0 \; , \quad D_T v^b \, + \, w^a_b v^b \, = \, 0 \; , \quad D_T g_{ab} \, = \, G_{ab} \label{eq:decomposition}$$

we easily verify that (18) holds.

If (u^a) , (v^b) are two pseudoparallel fields of virtual vectors than $D_T(g_{ab}u^av^b) = G_{ab}u^av^b$ and for this season the scalar function $g_{ab}u^av^b$ is generally not a constant. Similarly, for two T-parallel fields of virtual vectors (u^a) , (v^b) the scalar function $g_{ab}u^av^b$ is not in a general case a constant.

Let (v^a) be a field of virtual vectors defined along the trajectory (2). If there exists the function k = k(T), $T \in J$ such that

(19)
$$D_T(kv^a) = 0$$
, $k \neq 0$, $g_{ab}v^av^b = \text{konst.} > 0$

holds for every $T \in J$, then we call (v^a) a δ -parallel field. The vector field (v^a) is the δ -parallel field if and only if

(20)
$$\delta_T v^a \equiv D_T v^a + \frac{1}{2} \frac{G_{bc} v^b v^c}{g_{bc} v^b v^c} v^a = 0$$

holds for every $T \in J$.

The proof is analogical to the case of the rheonomous anholonomic manifold in [1],

¹⁾ I.e. the scalar product of two W-parallel "displaced" vectors is invariant.

(20) and (22). It is easy to verify the unique existence of the parallel displacement and δ -geodesic for the usual initial conditions. Let us remark that in entirely similar way as in [1] we may define the notion of *H*-parallel displacement on the anholonomic rheonomous manifold $r - R_n^m(t)$.

If the rheonomous manifold $r - R_n^m(t)$ is stationary than at every point $G_{ab} = 0$ and the δ -parallel displacement of vectors along the given trajectory is mutually identical with the T-parallel, W-parallel and pseudoparallel displacement.

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