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*Archivum Mathematicum*, Vol. 39 (2003), No. 1, 11--25

Persistent URL: <http://dml.cz/dmlcz/107850>

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**A FUNCTIONAL MODEL FOR A FAMILY OF OPERATORS INDUCED BY LAGUERRE OPERATOR**

HATAMLEH RA'ED

ABSTRACT. The paper generalizes the instruction, suggested by B. Sz.-Nagy and C. Foias, for operatorfunction induced by the Cauchy problem

$$T_t : \begin{cases} th''(t) + (1-t)h'(t) + Ah(t) = 0 \\ h(0) = h_0(th')(0) = h_1 \end{cases}$$

A unitary dilatation for  $T_t$  is constructed in the present paper. then a translational model for the family  $T_t$  is presented using a model construction scheme, suggested by Zolotarev, V., [3]. Finally, we derive a discrete functional model of family  $T_t$  and operator  $A$  applying the Laguerre transform

$$f(x) \rightarrow \int_0^\infty f(x) P_n(x) e^{-x} dx$$

where  $P_n(x)$  are Laguerre polynomials [6, 7]. We show that the Laguerre transform is a straightening transform which transfers the family  $T_t$  (which is not semigroup) into discrete semigroup  $e^{-itn}$ .

## INTRODUCTION

Functional models for contraction semigroups  $Z_t = \exp(itA)$  and  $T^n$ , ( $t \geq 0$ ,  $n \in \mathbb{Z}^+$ ) have been constructed by B. Sz.-Nagy and C. Foias [2] at the beginning of 70-s. The bases of this method is a significant concept of dilatation of contraction semigroup. A spectral realization of the dilatation and subsequent narrowing upon the original space leads to a functional model of the contraction semigroup. As a result an operator  $A(T)$  in this case is realized by operators which carry out multiplication by independent variable in a specific functional space. The basis of the concept is the Fourier transform of space  $L^2$ .

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2000 *Mathematics Subject Classification*: Primary: 47D06, 47A40; Secondary: 47A50, 47A48, 42A50.

*Key words and phrases*: Laguerre operator, semigroup, Hilbert space, functional model.

Received November 15, 1999.

1. PRELIMINARY INFORMATION ON THE FUNCTIONAL  
MODEL IN A FOURIER REPRESENTATION

**1.1.** We recall [1] that operator collegation  $\Delta$ ,

$$(1) \quad \Delta = (A, H, \phi, E, \sigma)$$

is a collection of Hilbert spaces  $H$  and  $E$  and of linear operators  $A : H \rightarrow H$ ,  $\phi : H \rightarrow E$ ,  $\sigma : E \rightarrow E$  ( $\sigma^* = \sigma$ ) where the collegation condition holds:

$$(2) \quad A - A^* = i\phi^* \sigma \phi.$$

It is customary to associate with the collegation (1) an open system [1] which is defined by relations

$$(3) \quad \begin{cases} i \frac{d}{dt} h(t) + Ah(t) = \phi^* \sigma u(t); \\ h(0) = h_0, \quad (t \geq 0); \\ v(t) = u(t) - i\phi h(t) \end{cases}$$

where  $h(t)$ ,  $u(t)$ ,  $v(t)$  are vector functions from Hilbert spaces  $H$  and  $E$  respectively. An important role in the further construction of the model representation plays the conservation Law [1].

**Theorem 1.1.** *For the open system (3) associated with the collegation  $\Delta$  (1) the conservation Law holds*

$$(4) \quad \|h_0\|^2 + \int_0^T \langle \sigma u(\zeta), u(\zeta) \rangle d\zeta = \|h(T)\|^2 + \int_0^T \langle \sigma v(\zeta), v(\zeta) \rangle d\zeta$$

for any  $T$ ,  $0 \leq T \leq \infty$ .

If operator  $A$  is selfadjoint then  $\phi = 0$ ,  $\sigma = 0$ , and Cauchy problem (3) in induced by the semigroup

$$Z_t = \exp(itA), \quad \text{i.e.} \quad h(t) = Z_t h_0$$

and the conservation Law (4) yields  $Z_t$ .

**1.2.** Let us consider a contractive semigroup  $Z_t = \exp(itA)$  ( $t \geq 0$ ), which has a property  $\|Z_t h\| \leq \|h\|$  for all  $h \in H$ .

A unitary dilatation of contractive semigroup  $Z_t$  in  $H$  is said to be a unitary semigroup  $U_t$  in  $\mathcal{H}$  [2] such that the following relation holds:

$$(5) \quad \mathcal{H} \supseteq H; \quad P_H U_t|_H = Z_t \quad (t \geq 0)$$

where  $P_H$  is an orthoprojector on  $H$ . The dilatation  $U_t$  in  $H$  is said to be minimal if

$$(6) \quad \mathcal{H} = \text{span}\{U_t h; t \in \mathbb{R}, h \in H\}$$

where span in (6) denotes a closed linear span of the vectors  $U_t h$  for any  $t \in \mathbb{R}$  and any  $h \in H$ .

A significant role in the theory of dilatation of contractive semigroup  $Z_t$  plays the following Theorem 1.2.

**Theorem 1.2.** *Any contracting semigroup  $Z_t$  in  $H$  has a unitary dilatation  $U_t$  in  $H$ . Moreover the minimal dilatation  $U_t$  is defined up to isomorphism.*

We present a construction of the dilatation  $U_t$  according to the paper [3]. A contractibility of the semigroup  $Z_t$  means [2, 3] that  $A$  is dissipative, i.e.  $-i(A - A^*) \geq 0$ . Consequently including  $A$  into the collocation  $\Delta$  (1) we can assume that  $\sigma = I$ . Therefore the conservation law (4) has the form

$$(7) \quad \|h_0\|^2 + \int_0^T \|u(\zeta)\|^2 d\zeta = \|h(T)\|^2 + \int_0^T \|v(\zeta)\|^2 d\zeta$$

We defined [3] a dilatation space  $\mathcal{H}$ , which forms vector-functions  $f(\zeta) = (u_+(\zeta), h, u_-(\zeta))$  so that  $u_{\pm}(\zeta) \in E$  and  $\text{Supp } u_{\pm}(\zeta) \in \mathbb{R}_{\mp}$  for a finite norm

$$(8) \quad \|f\|^2 = \int_{-\infty}^0 \|u_+(\zeta)\|^2 d\zeta + \|h\|^2 + \int_0^{\infty} \|u_-(\zeta)\|^2 d\zeta < \infty.$$

We define a dilatation  $U_t$  in  $\mathcal{H}$  by the formula

$$(9) \quad (U_t f)(\zeta) = (u_+(t, \zeta), h_t, u_-(t, \zeta))$$

where  $u_-(t, \zeta) = P_{\mathbb{R}_+} u_-(\zeta + t)$ ;  $h_t = y_t(0)$ , and  $y_t(\zeta)$  is a solution of the Cauchy problem

$$\begin{cases} i \frac{d}{d\zeta} y_t(\zeta) + A y_t(\zeta) = \phi^* u_-(\zeta + t); \\ y_t(-t) = 0, \quad \zeta \in (-t, 0); \end{cases}$$

and at last  $u_+(t, \zeta) = u_+(t + \zeta) + P_{(-t, 0)} \{u_-(\zeta + t) - i\phi y_t(\zeta)\}$  where  $P_{\mathbb{R}_+}$  and  $P_{(-t, 0)}$  are operators of narrowing (projection operators at set  $\mathbb{R}_+$  and  $(-t, 0)$  respectively),  $t \geq 0$ .

It is not difficult to show that unitary of  $U_t$  (9) in  $\mathcal{H}$  is a consequence of the conservation law (1). By the dilatation construction  $U_t$  one can see that the space  $\mathcal{H}$  has the form

$$(10) \quad \mathcal{H} = D_+ \oplus H \oplus D_-$$

where the subspace  $D_+$  is found by vector-function of the form  $(u_+(\zeta), 0, 0) \in \mathcal{H}$  and the subspace  $D_-$  is formed by vector-function  $(0, 0, u_-(\zeta))$  from  $\mathcal{H}$ , respectively.

The subspaces  $D_{\pm}$  have the following properties:

$$(11) \quad \begin{aligned} U_t D_+ &\subseteq D_+ & (t \geq 0), \\ U_t D_- &\subseteq D_- & (t \leq 0). \end{aligned}$$

Thus  $D_+$  is outgoing subspace and  $D_-$  is incoming subspace in the sense of P. D. Lax and R. S. Phillips [4]. In accordance with the paper [3], we define a free unitary group  $V_t$  in the space  $L^2_{\mathbb{R}}(E)$ , which will act as

$$(12) \quad (V_t g)(\zeta) = g(\zeta + t)$$

and vector-function  $g(\zeta) \in E$ ,  $\zeta \in \mathbb{R}$  is such that

$$\int_{-\infty}^{\infty} \|g(\zeta)\|^2 d\zeta < \infty.$$

It is evidently that  $D_{\pm}$  after identification belongs to  $L^2_{\mathbb{R}}(E)$  also.

Wave operators  $W_{\pm}$  play a significant role in the scattering theory. They are defined [3, 4] as

$$(13) \quad W_{\pm} = s - \lim_{t \rightarrow \mp\infty} U_+ P_{D_{\pm}} V_{-t}$$

where  $P_{D_{\pm}}$  are orthoprojectors on subspaces  $D_{\pm}$ . The following theorem holds [3].

**Theorem 1.3.** *The wave operators  $W_{\pm}$  exist as strong limits (13) are isometries from  $L^2_{\mathbb{R}}(E)$  to  $\mathcal{H}$ , and the relations*

$$(14) \quad W_{\pm} V_t = U_t W_{\pm}, \quad (\forall t), \quad W_{\pm} P_{D_{\pm}} = P_{D_{\pm}}$$

are valid.

The scattering operator  $S$  is defined by the wave operator  $W_{\pm}$  in a conventional way [3, 4]:

$$(15) \quad S = W_+^* W_-.$$

From Theorem 1.3 there follows a proposition.

**Theorem 1.4.** *The operator  $S$  (15) is a contraction, i.e.  $\|S\| \leq 1$  and has the properties:*

$$(16) \quad \begin{aligned} SV_t &= V_t S; & SL^2_{\mathbb{R}_+} &\subseteq L^2_{\mathbb{R}_+}(E); \\ \overline{SL^2_{\mathbb{R}}(E)} &= L^2_{\mathbb{R}}(E) \end{aligned}$$

**1.3.** We recall that the collegation  $\Delta$  (1) is simple [1-3] if  $H = \text{span}\{A^n \phi^* g; n \in \mathbb{Z}_+, g \in E\}$ . Let us define the following subspaces in  $\mathcal{H}$ ,

$$\mathfrak{R}_{\pm} = \overline{W_{\pm} L^2_{\mathbb{R}}(E)}.$$

The following theorem gives a sufficient condition for the completeness of the wave operators  $W_{\pm}$ , [3].

**Theorem 1.5.** *If the collegation  $\Delta$  is simple then the relation  $\mathcal{H} = \text{span}\{f_+ + f_-; f_{\pm} \in \mathfrak{R}_{\pm}\}$  holds.*

Now we construct a translational model [3]. Let  $f_k(\zeta) \in L^2_{\mathbb{R}}(E)$ , ( $k = 1, 2$ ). We define a mapping

$$\begin{pmatrix} f_1(\zeta) \\ f_2(\zeta) \end{pmatrix} \rightarrow \Psi_p(\zeta) = W_- f_1(\zeta) + W_+ f_2(\zeta) \in \mathcal{H}.$$

Then using isometry of  $W_{\pm}$  and the form of operator  $S$  (15) it is not difficult to show that

$$(17) \quad \|\Psi_p(\zeta)\|^2 = \int_{-\infty}^{\infty} \left\langle \begin{bmatrix} I & S^* \\ S & I \end{bmatrix} \begin{pmatrix} f_1(\zeta) \\ f_2(\zeta) \end{pmatrix}, \begin{pmatrix} f_1(\zeta) \\ f_2(\zeta) \end{pmatrix} \right\rangle d\zeta,$$

Using Theorem 1.5 we may assert, that space  $H$  is isomorphic to the space  $L^2 \left( \begin{smallmatrix} 1 & S^* \\ S & 1 \end{smallmatrix} \right)$  which is formed by vector-functions  $f(\zeta) = \begin{pmatrix} f_1(\zeta) \\ f_2(\zeta) \end{pmatrix}$  for which the norm (17) is finite. By virtue of conditions (14) the dilatation  $U_t$  on  $\Psi_p$  will act as a shift. Therefore if  $f(\zeta) \in L^2 \left( \begin{smallmatrix} 1 & S^* \\ S & 1 \end{smallmatrix} \right)$  then the dilatation  $U_t$  is transformed into

$$(18) \quad \widehat{U}_t f(\zeta) = f(\zeta + t).$$

Applying again (14), one can easily deduce that the spaces  $D_{\pm}$  are realized now in the form

$$(19) \quad \widehat{D}_- = \begin{pmatrix} L^2_{\mathbb{R}_+}(E) \\ 0 \end{pmatrix}, \quad \widehat{D}_+ = \begin{pmatrix} 0 \\ L^2_{\mathbb{R}_-}(E) \end{pmatrix}.$$

Thus the initial space  $H$  acquires such model form

$$(20) \quad \begin{aligned} \widehat{H}_p &= L^2 \left( \begin{smallmatrix} 1 & S^* \\ S & 1 \end{smallmatrix} \right) \ominus \begin{pmatrix} L^2_{\mathbb{R}_+}(E) \\ L^2_{\mathbb{R}_-}(E) \end{pmatrix} \\ &= f = \left( \begin{pmatrix} f_1 \\ f_2 \end{pmatrix} \in L^2 \left( \begin{smallmatrix} 1 & S^* \\ S & 1 \end{smallmatrix} \right); \begin{matrix} f_1 + S^* f_2 \in L^2_{\mathbb{R}_-}(E) \\ S f_1 + f_2 \in L^2_{\mathbb{R}_+}(E) \end{matrix} \right) \end{aligned}$$

and in the virtue of the dilatation the action of semigroup  $Z_t$  is transformed to the shift semigroup

$$(21) \quad \widehat{Z}f(\zeta) = P_{\widehat{H}_p} f(\zeta + t)$$

where  $f(\zeta) \in \widehat{H}_p$  (20). Thus the following theorem is proved.

**Theorem 1.6.** *A minimal unitary dilatation  $U_t$  in  $\mathcal{H}$  of the contraction semigroup  $Z_t = \exp(itA)$  in  $H$ , where  $A$  is dissipative operator of a simple collocation  $\Delta$  is unitary equivalent to a translation group  $\widehat{U}_t$  (18) in the space  $L^2 \left( \begin{smallmatrix} 1 & S^* \\ S & 1 \end{smallmatrix} \right)$ , and the contraction semigroup  $Z_t$  is unitary equivalent to the shift semigroup  $\widehat{Z}_t$  (21) in the space  $\widehat{H}_p$  respectively.*

The Fourier transform by formula

$$(22) \quad \widetilde{f}(\lambda) = \int_{-\infty}^{\infty} f(\zeta) e^{-i\lambda\zeta} d\zeta$$

in the virtue of Plancherel theorem [2, 3] is a unitary operator in  $L^2_{\mathbb{R}}(E)$ . By the virtue of Wiener-Paley theorem

$$\tilde{L}^2_{\mathbb{R}_+}(E) = H^2_-(E); \quad \tilde{L}^2_{\mathbb{R}_-}(E) = H^2_+(E)$$

where  $H^2_{\pm}(E)$  are Hardy spaces of  $E$ -value function from  $L^2_{\mathbb{R}}(E)$  which are holomorphically continued into lower (upper) half-plane. Let us apply the Fourier transform (22) to translational model (18) – (21) and take advantage of the following Theorem 1.7

**Theorem 1.7.** *The Fourier transform of the scattering operator  $S$  (15) transfers the operator  $S$  into operator performing multiplication by characteristic function*

$$(23) \quad \begin{aligned} S_{\Delta}(\lambda) &= I - \phi(A - \lambda I)^{-1} \phi^*, \quad i.e. \\ (\tilde{S}f)(\lambda) &= S_{\Delta}(\lambda) \tilde{f}(\lambda). \end{aligned}$$

As it is known  $\tilde{f}(\lambda + t) = e^{i\lambda t} \tilde{f}(\lambda)$ , therefore we derive such functional model.

**Theorem 1.8.** *A minimal unitary dilatation  $U_t$  in  $H$  of the contraction semigroup  $Z_t = \exp(itA)$  in  $H$ , where  $A$  is dissipative operator of a simple collegation  $\Delta$  is unitary equivalent to the group*

$$(24) \quad \tilde{U}_t f(\lambda) = e^{i\lambda t} f(\lambda)$$

where  $f(\lambda) \in L^2 \left( \begin{array}{cc} I & S_{\Delta}^*(\lambda) \\ S_{\Delta}(\lambda) & I \end{array} \right)$  and contraction semigroup  $Z_t$  is unitary equivalent to semigroup  $\tilde{Z}_t f(\lambda) = P_{\tilde{H}_p} e^{i\lambda t} f(\lambda)$ , where  $f(\lambda)$  belongs to the space

$$\tilde{H}_p = \left\{ f = \begin{pmatrix} f_1 \\ f_2 \end{pmatrix}(\lambda) \in \begin{pmatrix} I & S_{\Delta}^*(\lambda) \\ S_{\Delta}(\lambda) & I \end{pmatrix}; f_1 + S_{\Delta}^*(\lambda) f_2 \in H^2_+(E) \right\} \\ \left. ; S_{\Delta}(\lambda) f_1 + f_2 \in H^2_-(E) \right\}$$

Here the main operator  $\tilde{A}$  in  $\tilde{H}_p$  act as multiplication operator by independent variable

$$(26) \quad \tilde{A}f(\lambda) = P_{\tilde{H}_p} \lambda f(\lambda), \quad f(\lambda) \in \tilde{H}_p.$$

In the next section we will generalize this construction on the case of the Laguerre transform.

## 2. A FUNCTIONAL MODEL FOR THE LAGUERRE REPRESENTATION

**2.1.** Let us consider a differential operator

$$(27) \quad \ell = t \frac{d^2}{dt^2} + (1-t) \frac{d}{dt}$$

in what follows called the Laguerre operator; it acts on functions form  $C^2 = (\mathbb{R}_+)$ . We denote by  $L^2_{\mathbb{R}_+}(e^{-t} dt)$  the following space:

$$(28) \quad L^2_{\mathbb{R}_+}(e^{it} dt) = \left\{ f(t), t \in \mathbb{R}_+; \int_0^{\infty} |f(t)|^2 e^{-t} dt < \infty \right\}$$

**Proposition 2.1.** *An operator  $\ell$  is symmetric in the space  $L^2_{\mathbb{R}_+}(e^{-t} dt)$  under the self-adjoint boundary conditions, i.e.  $\langle \ell x, y \rangle = \langle y, \ell y \rangle$  for all  $x, y \in \mathbb{C}^2(\mathbb{R}_+)$  such that  $tx(t)|_{t=0} = 0$ ,  $ty(t)|_{t=0} = 0$  and  $ty'(t)|_{t=0} < \infty$ ,  $tx'(t)|_{t=0} < \infty$ .*

**Proof.** We calculate

$$\begin{aligned} \langle \ell x, y \rangle - \langle x, \ell y \rangle &= \int_0^\infty \{ (tx'' + (1-t)x')\bar{y} - x(t\bar{y}'' + (1-t)\bar{y}') \} e^{-t} dt \\ &= \int_0^\infty \{ te^{-t}(x'\bar{y} - \bar{y}'x) \}' dt = \{ te^{-t}(x'\bar{y} - \bar{y}'x) \} \Big|_0^\infty = 0 \end{aligned}$$

by virtue of the boundary conditions.  $\square$

Let us consider now an open system of special form, generated by the Laguerre operator (27) and corresponding to the collegation  $\Delta$  (1):

$$(29) \quad \begin{cases} \ell h(t) + Ah(t) = \phi^* \sigma u(t); \\ h(0) = h_0(th')(0) = h_1; \\ v(t) = u(t) - i\phi h(t). \end{cases}$$

The following assertion is valid, similar to Theorem (1.1).

**Theorem 2.1.** *For the open system (29) associated with collegation  $\Delta$  the law of conservation of energy is valid, i.e.*

$$(30) \quad \begin{aligned} &\int_0^T \langle \sigma u(\zeta), u(\zeta) \rangle e^{-\zeta} d\zeta + \langle \widehat{I}h_0, \widehat{h}_0 \rangle \\ &= \int_0^T \langle \sigma v(\zeta), v(\zeta) \rangle e^{-\zeta} d\zeta + \langle \widehat{I}h_T, \widehat{h}_T \rangle \end{aligned}$$

where  $I = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$  and  $h_0 = \begin{pmatrix} h_0 \\ h_t \end{pmatrix}$ ,  $h_T = \begin{pmatrix} h(T) \\ e^{-T}Th'(T) \end{pmatrix}$  for any finite  $T > 0$ .

**Proof.** We calculate

$$\begin{aligned} \langle \ell h, h \rangle - \langle h, \ell h \rangle &= \langle \phi^* \sigma u - Ah, h \rangle - \langle h, \psi^* \sigma u - Ah \rangle \\ &= \langle \sigma u, \frac{u-v}{i} \rangle - \langle \frac{u-v}{i}, \sigma u \rangle - \langle (A - A^*)h, h \rangle \\ &= i \langle \sigma u, u-v \rangle + i \langle u-v, \sigma u \rangle - i \langle \phi^* \sigma \phi h, h \rangle \\ &= i \langle \sigma u, u-v \rangle + i \langle u-v, \sigma u \rangle - i \langle \sigma(u-v), u-v \rangle \\ &= i \langle \sigma u, u \rangle - i \langle \sigma v, v \rangle. \end{aligned}$$

Now we integrate the derived equality:

$$\begin{aligned} &\int_0^T \langle \sigma v, v \rangle e^{-t} dt - \int_0^T \langle \sigma u, u \rangle e^{-t} dt \\ &= i \int_0^T [\langle \ell h, h \rangle - \langle h, \ell h \rangle] e^{-t} dt \\ &= i \{ e^{-t} t [\langle h'', h \rangle - \langle h, h' \rangle] \} \Big|_0^T \\ &= \langle \widehat{I}h_0, \widehat{h}_0 \rangle - \langle \widehat{I}h_T, \widehat{h}_T \rangle \end{aligned}$$



which proves our assertion.  $\square$

**2.2.** Let us make use of the energy conservation law (30) to construct a dilatation for operator  $T_t$  generated by the Cauchy problem

$$(31) \quad \begin{cases} \ell h(t) + Ah(t) = 0; \\ h(0) = h_0; (th')(0) = h_1; \end{cases}$$

where  $T_t(h_0, h_1) = (h(t), th'(t))$ . We will call an unitary operator-function  $U_t$  in  $\mathcal{H}$  a dilatation of family  $T_t$  in  $H$ , if  $\mathcal{H} \supseteq H$ ,  $T_t = P_H U_t|_H$ .

Here we do not suppose that  $T_t$  and  $U_t$  is semigroup. Moreover, the unitary property of  $U_t$  may hold not necessarily in Hilbert metric but in indefinite one. The following analog of Theorem 1.2 is valid.

**Theorem 2.2.** *The operator-function  $T_t$  generated by the Cauchy problem (31) with dissipative operator  $A$  of collegation  $\Delta$  (1) (i.e.  $\sigma = I$ ) possesses the unitary (in indefinite metric) dilatation  $U_t$ , where the minimal dilatation is determined up to isomorphism.*

**Proof.** To prove the theorem we bring a construction of dilatation  $U_t$  by analog with (8), (9).

Let us consider a Hilbert space

$$(32) \quad \begin{aligned} \mathcal{H} = \{ & f = (u(\zeta), \widehat{h}, v(\zeta)); u(\zeta), v(\zeta) \in E, \text{supp } v \in \mathbb{R}_-, \text{supp } u \in \mathbb{R}_+, \\ & \widehat{h} = \begin{pmatrix} h_0 \\ h_1 \end{pmatrix}, h_k \in H; \|f\|^2 = \int_{-\infty}^0 \|v(\zeta)\|^2 e^{-\zeta} d\zeta + \|\widehat{h}\|^2 \\ & + \int_0^{\infty} \|u(\zeta)\|^2 e^{-\zeta} d\zeta < \infty \}. \end{aligned}$$

We set indefinite metric  $\mathcal{H}$

$$(33) \quad \langle f \rangle_I^2 = \int_{-\infty}^0 \|v(\zeta)\|^2 e^{-\zeta} d\zeta + \langle I\widehat{h}, \widehat{h} \rangle + \int_0^{\infty} \|u(\zeta)\|^2 e^{-\zeta} d\zeta$$

where  $I$  has the form indicated in Theorem 2.1.

We construct the dilatation  $U_t$  in  $\mathcal{H}$ ,

$$(34) \quad U_t f = f_t(u(t, \zeta), \widehat{h}_t, v(t, \zeta)).$$

Let us consider further the Cauchy problem

$$(35) \quad \begin{cases} (i \frac{\partial}{\partial t} + \ell_\zeta) \widehat{u}(t, \zeta) = 0; \\ \widehat{u}(0, \zeta) = u(\zeta); \zeta \in \mathbb{R}_+; \end{cases}$$

where  $\ell_\zeta$  is operator  $\ell$  (27) with respect to  $\zeta$ .

Solution of the problem is easily obtained. In fact, let

$$\widehat{u}(t, \zeta) = \sum_{n \in \mathbb{Z}_+} e^{-itn} C_n g_n(\zeta)$$

where  $g_n(\zeta)$  are the Laguerre polynomials [5] which are the solutions of equation  $\ell_\zeta g_n(\zeta) + n g_n(\zeta) = 0$  and have the form

$$g_n(\zeta) = \frac{1}{n!} e^\zeta \frac{d^n}{d\zeta^n} (\zeta e^{-\zeta})$$

and make a complete system of orthogonal polynomials in  $L^2_{\mathbb{R}_+}(e^{-\zeta} d\zeta)$ . The coefficients  $C_n$  are obtained from the initial condition  $\sum C_n g_n(\zeta) = u(\zeta)$ .

Therefore  $\widehat{u}(t, \zeta)$  possesses the property  $\text{supp } \widehat{u}(t, \zeta) = \text{supp } \widehat{u}(\zeta) \subseteq \mathbb{R}_+$ . Now we determine  $u(t, \zeta)$  in (34) by the formula

$$(36) \quad u(t, \zeta) = P_{\mathbb{R}_+} \widehat{u}(t, \zeta + t) e^{-\frac{t}{2}}.$$

To set  $\widehat{h}_t$  (34), we consider the following Cauchy problem

$$(37) \quad \begin{cases} \ell_\zeta y(\zeta) + Ay(\zeta) = \phi^* \widehat{u}(t, \zeta + t) e^{-\frac{t}{2}}; & \zeta \in (-t, 0); \\ y(-t) = h_0; \\ (-t)e^t y(-t) = h_1; \end{cases}$$

and put  $\widehat{h}_t = \begin{pmatrix} y(0) \\ (ty')(0) \end{pmatrix}$ .

Finally, to set  $v(t, \zeta)$  (34) we consider the similar equation

$$(38) \quad \begin{cases} (i \frac{\partial}{\partial t} + \ell_\zeta) \widehat{v}(t, \zeta) = 0; \\ \widehat{v}(0, \zeta) = v(\zeta); & \zeta \in \mathbb{R}_-; \end{cases}$$

and put  $v(t, \zeta) = e^{-\frac{t}{2}} \widehat{v}(t, \zeta + t) + P_{\mathbb{R}_-} \{ \widehat{u}(t, \zeta + t) e^{-\frac{t}{2}} - i\phi y(\zeta) \}$ . We show that  $U_t$  (34) has property of isometry in the metric (33). To this end we calculate,

$$\begin{aligned} \langle f_t \rangle_I^2 &= \int_{-\infty}^0 \|v(t, \zeta)\|^2 e^{-\zeta} d\zeta + \langle I \widehat{h}_t, \widehat{h}_t \rangle + \int_0^\infty \|u(t, \zeta)\|^2 e^{-\zeta} d\zeta \\ &= \int_{-\infty}^{-t} \|\widehat{v}(t, \zeta + t)\|^2 e^{-\zeta - t} d\zeta + \int_{-t}^0 \|\widehat{u}(t, \zeta + t) e^{-\frac{t}{2}} - i\phi y(\zeta)\|^2 e^\zeta d\zeta \\ &\quad + \langle I \widehat{h}_t, \widehat{h}_t \rangle + \int_0^\infty \|u(t, \zeta + t)\|^2 e^{-\zeta - t} d\zeta \\ &= \int_{-\infty}^{-t} \|\widehat{v}(t, \zeta + t)\|^2 e^{-\zeta - t} d\zeta + \langle I \widehat{h}_0, \widehat{h}_0 \rangle + \int_{-t}^\infty \|\widehat{u}(t, \zeta + t)\|^2 e^{-\zeta - t} d\zeta \\ &= \int_{-\infty}^0 \|\widehat{v}(t, \zeta)\|^2 e^{-\zeta} d\zeta + \langle I \widehat{h}_0, \widehat{h}_0 \rangle + \int_0^\infty \|\widehat{u}(t, \zeta)\|^2 e^{-\zeta} d\zeta \\ &= \langle f \rangle_I^2 \end{aligned}$$

In this calculation we have made use of the conservation law (30) and of the fact that norms of solutions of Cauchy problems  $\widehat{u}(t, \zeta)$ ,  $\widehat{v}(t, \zeta)$  (35) and (38) coincide with norms of initial data  $u(\zeta)$  and  $v(\zeta)$  in the spaces  $L_{\mathbb{R}_+}^2(e^{-t} dt)$  and  $L_{\mathbb{R}_-}^2(e^{-t} dt)$  by virtue of selfadjointness of operators  $\ell_\zeta$  in the spaces.

In order to prove that  $U_t$  has the property of being unitary, it is necessary to ascertain that from  $U_t^* f = 0$  implies  $f = 0$ . It is easy to show that  $U_t^*$  will act by the formula

$$(39) \quad U_t^* f = (u(t, \zeta), \widehat{h}_t, v(t, \zeta)).$$

Here  $v(t, \zeta) = P_{\mathbb{R}_-} \widehat{v}(t, \zeta - t)e^{\frac{t}{2}}$  where  $\widehat{v}(t, \zeta)$  is a solution of problem (38).

In order to obtain  $\widehat{h}_t$ , it is necessary to consider dual to (37) problem

$$(40) \quad \begin{cases} \ell_\zeta y(\zeta) + A^* y(\zeta) = \phi^* \widehat{v}(\zeta, \zeta - t)e^{\frac{t}{2}}; \\ y(t) = h_0; \\ e^{-t} t y'(t) = h_1; \end{cases}$$

and put  $\widehat{h}_t = \begin{pmatrix} y(0) \\ (ty')(0) \end{pmatrix}$ . Finally,

$$u(t, \zeta) = \widehat{u}(t, \zeta - t)e^{\frac{t}{2}} + P_{\mathbb{R}_+} \{ \widehat{v}(t, \zeta - t)e^{\frac{t}{2}} + i\phi y(\zeta) \},$$

where  $\widehat{u}(t, \zeta)$  is the solution of Cauchy problem (35).

Thus let  $U_t^* f = 0$ , then  $\widehat{u}(t, \zeta) = 0$  and so  $\widehat{u}(t, \zeta) = 0$  and  $\widehat{v}(t, \zeta - t)e^{\frac{t}{2}} + i\phi y(\zeta) = 0$  therefore  $u(\zeta) \equiv 0$ . Now, by substituting  $\widehat{v}(t, \zeta - t) = -i\phi y(\zeta)e^{-\frac{t}{2}}$  in (40) we obtain a homogeneous equation

$$\ell_\zeta y + A^* y + i\phi^* \phi y = 0$$

with zero condition in the origin  $\widehat{h}_t = 0$ . By virtue of uniqueness of Cauchy problem solution, this yields that  $y(\zeta) \equiv 0$ , therefore  $\widehat{v}(t, \zeta - t) = 0$  on interval  $(0, t)$ . Accounting that  $\widehat{v}(t, \zeta - t) = 0$  with  $(-\infty, 0)$ , finally we conclude that  $v(\zeta) = 0$ . Thus  $f = 0$ . This proves the property of being unitary for  $U_t$  (34) and completes the proof of the theorem.  $\square$

**2.3.** Let us pass to constructing wave operators. To this end we define a “free” group by analogy with (38)

$$(41) \quad V_t g(\zeta) = g(t, \zeta),$$

where  $g(t, \zeta)$  is a solution of Cauchy problem

$$(42) \quad \begin{cases} (i\frac{\partial}{\partial t} + \ell_\zeta) g(t, \zeta) = 0; \\ g(0, \zeta) = g(\zeta) \in L_{\mathbb{R}}^2(e^{-\zeta} d\zeta). \end{cases}$$

It is evident that  $V_t$  (41) is unitary. Now we define the operators

$$(43) \quad \begin{aligned} W_- &= s - \lim_{t \rightarrow +\infty} U_t P_{\mathbb{R}_+} V_{-t}, \\ W_+ &= s - \lim_{t \rightarrow -\infty} U_t^* P_{\mathbb{R}_-} V_{-t}^*. \end{aligned}$$

By analogy with Theorem 1.3 we have

**Theorem 2.3.** *The wave operators  $W_{\pm}$  exist as strong limits (43), are isometries from  $L^2_{\mathbb{R}}(e^{-\zeta} d\zeta)$  to  $\mathcal{H}$ , and the following relations are valid:*

$$(44) \quad \begin{aligned} U_t W_- &= W_- V_t, U_t^* W_+ = W_+ V_t^*, \quad (t \geq 0) \\ W_{\pm} P_{\mathbb{R}_{\mp}} &= P_{\mathbb{R}_{\mp}} \end{aligned}$$

**Proof.** We prove the assertion of the theorem for  $W_-$  (for  $W_+$  the proof is similar). The main matter of the theorem consists of existence proof of  $W_-$  since the relation (44) is proved by analogy with arguments given in Section 1; see [2, 3]. Let

$$f_t = U_t P_{\mathbb{R}_+} V_{-t} g = (v(t, \zeta), h_t, u(t, \zeta))$$

then  $u(t, \zeta) = P_{\mathbb{R}_+} g(\zeta)$ . We consider the Cauchy problem

$$(45) \quad \begin{cases} \ell_{\zeta} y(\zeta) + Ay(\zeta) = \phi^* g(\zeta); \\ y(-t) = 0; y'(-t) = 0, \quad \zeta \in (-t, 0). \end{cases}$$

Then  $\widehat{h}_t = \begin{pmatrix} y(0) \\ (ty')(0) \end{pmatrix}$ .

We denote by  $K(\zeta, \eta)$  a Cauchy function of the problem (45) (i.e.  $K(\zeta, \zeta) = 0$ ,  $K'(\zeta, \zeta) = I$ ), then a solution  $y(\zeta)$  of (45) has the form

$$y_t(\zeta) = \int_{-t}^{\zeta} K(\zeta, \eta) \phi^* g(\eta) d\eta.$$

Therefore  $V(t, \zeta)$  has the form

$$V(t, \zeta) = P_{(-t, 0)} \{g(\zeta) - i\phi y(\zeta)\}.$$

Thus,

$$f_t = \left( P_{(-t, 0)} \{g(\zeta) - i\phi \int_{-t}^0 K(\zeta, \eta) \phi^* g(\eta) d\eta\}, \begin{pmatrix} \int_{-t}^0 K(0, \eta) \phi^* g(\eta) d\eta \\ \int_{-t}^0 K'(0, \eta) \phi^* g(\eta) d\eta \end{pmatrix}, P_{\mathbb{R}_+} g(\zeta) \right).$$

We show that  $f_t$  is a Cauchy sequence, i.e.  $\|f_{t+\Delta} - f_t\|^2 \rightarrow 0$  as  $t \rightarrow \infty$ . Since

$$(46) \quad \|f_{t+\Delta} - f_t\|^2 = \int_{-\infty}^0 \|v_t(t + \Delta, \zeta) - v(t, \zeta)\|^2 e^{-\zeta} d\zeta + \|\widehat{h}_{t+\Delta} - \widehat{h}_t\|^2.$$

It is sufficient to show that each summand approaches to zero as  $t \rightarrow \infty$ . We show that  $\|\widehat{h}_{t+\Delta} - \widehat{h}_t\| \rightarrow 0$  when  $t \rightarrow \infty$  and we will prove this property component by component. It is obvious that

$$\begin{aligned} \|\widehat{h}_{t+\Delta} - \widehat{h}_t\|^2 &= \left\| \int_{(-t-\Delta)}^{-t} K(0, \eta) \phi^* g(\eta) d\eta \right\|^2 \\ &\leq \int_{-t-\Delta}^{-t} \|K(0, \eta)\|^2 e^{\eta} d\eta \cdot \int_{-t-\Delta}^{-t} e^{-\eta} \|\phi^*\|^2 \|g(\eta)\|^2 d\eta \end{aligned}$$

and since the function  $K(0, \eta)e^\eta$  is bounded (see [6, 7]), we obtain that

$$\|h_{t+\Delta} - h_t\|^2 \leq \Delta C \|\phi^*\|^2 \int_{-t-\Delta}^{-t} \|g(\eta)\|^2 e^{-\eta} d\eta \rightarrow 0 \quad \text{as } t \rightarrow \infty$$

since  $g(\eta) \in L_{\mathbb{R}}^2(e^{-\eta} d\eta)$ .

The convergence of second components  $\widehat{h}_{t+\Delta} - \widehat{h}_t$  to zero is proved in a similar way. We show that the first summand in (46) approaches to zero too.

In fact,

$$\begin{aligned} A &= \int_{-\infty}^0 \|P_{(-t-\Delta, -t)}g(\zeta) - iP_{(-t-\Delta, 0)}\phi \int_{-t-\Delta}^{\zeta} K(\zeta, \eta)\phi^*g(\eta) d\eta \\ &\quad + i \int_{-t}^{\zeta} \phi K(\zeta, \eta)\phi^*g(\eta) d\eta\|^2 e^{-\zeta} d\zeta \\ &= \int_{-t-\Delta}^{-t} \|g(\zeta)\|^2 e^{-\zeta} d\zeta + \int_{-\infty}^0 \|P_{(-t-\Delta, 0)}\phi y_{t+\Delta}(\zeta) - P_{(-t, 0)}\phi y_t(\zeta)\|^2 e^{-\zeta} d\zeta \\ &\quad + 2\text{Im} \int_{-t-\Delta}^{-t} \langle g(\zeta), P_{(-t-\Delta, 0)}\phi y(\zeta) - P_{(-t, 0)}\phi y(\zeta) \rangle e^{-\zeta} d\zeta \end{aligned}$$

It is obvious that the first and third summands in the given sum approaches to zero as  $t \rightarrow \infty$  because  $g(\zeta) \in L_{\mathbb{R}}^2(e^{-\zeta} d\zeta)$ . We evaluate the second summand:

$$\begin{aligned} B &= \int_{-\infty}^0 \|P_{(-t-\Delta, 0)}\phi y_{t+\Delta}(\zeta) - P_{(-t, 0)}\phi y_t(\zeta)\|^2 e^{-\zeta} d\zeta \\ &= \int_{-\infty}^0 \langle \phi \Delta y, \phi \Delta y \rangle e^{-\zeta} d\zeta, \end{aligned}$$

where

$$\Delta y = P_{(-t-\Delta, 0)}y_{t+\Delta}(\zeta) - P_{(-t, 0)}y_t(\zeta).$$

Then

$$\begin{aligned} A &= \int_{-\infty}^0 \langle \phi^* \phi \Delta y, \Delta y \rangle e^{-\zeta} d\zeta = \int_{-\infty}^0 \left\langle \frac{A - A^*}{i} \Delta y, \Delta y \right\rangle e^{-\zeta} d\zeta \\ &= 2\text{Im} \int_{-\infty}^0 \langle \phi^* g - \ell \Delta y, \Delta y \rangle e^{-\zeta} d\zeta \\ &= 2\text{Im} \int_{-\infty}^0 \langle \phi^* g, \Delta y \rangle e^{-\zeta} d\zeta + 2\text{Im} \int_{-\infty}^0 \langle \ell \Delta y, \Delta y \rangle e^{-\zeta} d\zeta \end{aligned}$$

the first summand approaches to zero again on account of  $g(\zeta) \in L_{\mathbb{R}}^2(e^{-\zeta} d\zeta)$ , and the second one yields after integration by parts

$$\|\zeta e^{-\zeta} \Delta y\|_{\zeta=0} \rightarrow 0 \quad (t \rightarrow \infty)$$

since  $\Delta \widehat{h}_t \rightarrow 0$ . The theorem is proved.  $\square$

As before, we define the operator  $S$  by the formula (15). Then the following theorem holds.

**Theorem 2.4.** *The operator  $S$  (15) is a contraction from  $L^2_{\mathbb{R}}(e^{-\zeta} d\zeta)$  to  $L^2_{\mathbb{R}}(e^{-\zeta} d\zeta)$  and possesses the following properties:*

$$SV_t = V_t S; \quad SL^2_{\mathbb{R}_+}(e^{-\zeta} d\zeta) \subset L^2_{\mathbb{R}_+}(e^{-\zeta} d\zeta);$$

$$\overline{SL^2_{\mathbb{R}}(e^{-\zeta} d\zeta)} = L^2_{\mathbb{R}}(e^{-\zeta} d\zeta).$$

**2.4.** Further we suppose that the collegation  $\Delta$  (1) is simple and as in subsection 1.3 we set a mapping

$$\Psi_p(\zeta = W_- f_1(\zeta)) + W_+ f_2(\zeta)$$

from  $L^2_{\mathbb{R}}(e^{-\zeta} d\zeta) + L^2_{\mathbb{R}}(e^{-\zeta} d\zeta)$  to  $\mathcal{H}$ . It is obvious that

$$\Psi_p(\zeta) \in L^2 \left( \begin{pmatrix} I & S^* \\ S & I \end{pmatrix}, e^{-\zeta} d\zeta \right)$$

Action of dilatation in this space again reduces to a translation

$$(47) \quad \widehat{U}_t f(\zeta) = f(\zeta + t),$$

since

$$\begin{aligned} U_t \Psi_p(\zeta) &= W_- f_1(\zeta + t) + U_t W_+ f_2(\zeta) \\ &= W_- f_1(\zeta + t) + U_t W_+ V_t^* V_t f_2(\zeta) \\ &= W_- f_1(\zeta + t) + U_t U_t^* W_+ V_t f_2(\zeta) = \Psi_P(\zeta + t). \end{aligned}$$

As earlier, it is obvious that

$$D_- = \begin{pmatrix} L^2_{\mathbb{R}_+}(e^{-\zeta} d\zeta) \\ 0 \end{pmatrix}, \quad D_+ = \begin{pmatrix} 0 \\ L^2_{\mathbb{R}_-}(e^{-\zeta} d\zeta) \end{pmatrix}$$

and the model space  $H_p$  has the form

$$(48) \quad H_p = L^2 \left( \begin{pmatrix} 1 & S^* \\ S & 1 \end{pmatrix} e^{-\zeta} d\zeta \right) \ominus \begin{pmatrix} L^2_{\mathbb{R}_+}(e^{-\zeta} d\zeta) \\ L^2_{\mathbb{R}_-}(e^{-\zeta} d\zeta) \end{pmatrix}$$

and in addition  $T_t$  passes to shift semigroup

$$(49) \quad \widehat{T}_t f(\zeta) = f(\zeta + t).$$

Now we consider a Laguerre transform

$$(50) \quad L_n = \int_0^\infty e^{-x} P_n(x) f(x) dx$$

where  $P_n(x) = \frac{1}{n!} e^{-x} \frac{d^n}{dx^n} (x e^{-x})$  are a Laguerre polynomials, and  $f(x) \in L^2_{\mathbb{R}_+}(e^{-x} dx)$ . The transform (50) ascertains isomorphism between  $L^2_{\mathbb{R}_+}(e^{-x} dx)$  and  $\ell^2$ .

We extend the Laguerre transform (50) on  $\mathbb{R}_-$  in a symmetric way. Then an image of this map yields a space  $\ell^2$ . Let  $\ell^2_{\mathbb{Z}} = \ell^2_- + \ell^2_+$  is a space of square summable two-sided sequences. Just as for the case of Fourier transform (see Theorem 1.7 in Section 1) a theorem the proof of which repeats the reasonings brought out in [3] holds.

**Theorem 2.5.** *The Laguerre transform of scattering operator  $S$  transfers the operator  $S$  into an operator of multiplication by a characteristic function  $S_\Delta(n) = I - i\phi(A - nI)^{-1}\phi^*$ ,  $n \in \mathbb{Z}$ , i.e.*

$$(51) \quad L_n(Sg) = S_\Delta(n)g_n$$

where  $g_n = L_n(g)$ .

After realizing the Laguerre transform, the space  $L^2\left(\left(\begin{smallmatrix} I & S^* \\ S & I \end{smallmatrix}\right) e^{-\zeta} d\zeta\right)$  passes into the space  $\ell_{\mathbb{Z}}^2\left(\begin{smallmatrix} I & S_\Delta^*(n) \\ S_\Delta(n) & I \end{smallmatrix}\right)$  and dilatation  $\widehat{U}_t$  (47) is converted into

$$(52) \quad \widehat{U}_t(n)f_n = e^{-itn}f_n.$$

Supspaces  $D_\pm$  will have the form

$$D_- = \begin{pmatrix} \ell_-^2 \\ 0 \end{pmatrix}, \quad D_+ = \begin{pmatrix} 0 \\ \ell_+^2 \end{pmatrix}.$$

Therefore  $H_p$  is converted to the form

$$(53) \quad \widetilde{H}_p = \left\{ f_n = \begin{pmatrix} f_n^1 \\ f_n^2 \end{pmatrix} \in \ell_{\mathbb{Z}}^2\left(\begin{smallmatrix} I & S_\Delta^*(n) \\ S_\Delta(n) & I \end{smallmatrix}\right); \begin{matrix} f_n^1 + S_\Delta^*(n)f_n^2 \in \ell_+^2 \\ S_\Delta(n)f_n^1 + f_n^2 \in \ell_-^2 \end{matrix} \right\}$$

and a “semigroup”  $T_t$  will have the form

$$(54) \quad \widetilde{T}_t(n)f_n = P_{\widetilde{H}_p} e^{-itn}f_n.$$

Thus the following theorem is proved.

**Theorem 2.6.** *The minimal unitary dilatation  $U_t$  (34) in  $\mathcal{H}$  (32) of the family of operators  $T_t$  (31) with a scattering operator  $A$  of collocation  $\Delta$  (1) is unitary equivalent to  $\widehat{U}_t(n)$  (52) in the space  $\ell_{\mathbb{Z}}^2\left(\begin{smallmatrix} I & S_\Delta^*(n) \\ S_\Delta(n) & I \end{smallmatrix}\right)$ , and the family  $T_t$  (31) is unitary equivalent to  $\widetilde{T}_t(n)$  (54) in the space  $\widetilde{H}_p$ .*

**Acknowledgement.** The author is grateful to V. A. Zolotarev for setting of the problem and his permanent interest in the work.

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