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ON GENERALIZED LOCALIZABILITY

VÁCLAV ALDA

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The localizability of a physical system S in a space X is defined as a homomorphism h of the lattice $B(X)$ of the Borel sets in the space X into the lattice \mathcal{E} of yes – no experiments which can be performed on S . In each state $s \in \mathcal{S}$ on \mathcal{E} and $A \in B(X)$, $s(h(A))$ is the probability of finding S in A . In [1] Jauch and Piron introduced generalized localizability; the condition that $h: B(X) \rightarrow \mathcal{E}$ is a homomorphism is weakened. They demand

- (1) $h(\emptyset) = 0, \quad h(X) = I,$
- (2) $A_1 \cap A_2 = \emptyset \Rightarrow h(A_1) \perp h(A_2),$
- (3) $h(A_1 \cap A_2) = h(A_1) \wedge h(A_2),$

but they do not demand $h(A_1 \cup A_2) = h(A_1) \vee h(A_2)$ for $A_1 \cap A_2 = \emptyset$. In their paper [1] \mathcal{E} is the lattice of projections in a Hilbert space \mathfrak{H} .

A transformation $h: B(X) \rightarrow \mathcal{E}$ satisfying (1)–(3) can be constructed in a simple manner. We take a homomorphism $h': B(X) \rightarrow \mathcal{E}$, a projection P non-commuting with all $h'(A)$, $A \in B(X)$, and then

$$h: h(A) = P \wedge h'(A)$$

satisfies (1)–(3) in $P\mathfrak{H}$.

This is the situation which is described in a theorem due to Neumark [2]: For every positive operator valued measure T in a Hilbert space \mathfrak{H} one can construct an extension $\mathfrak{H}' \supset \mathfrak{H}$ and a projection valued measure T' so that

$$T(A) = P T'(A) P$$

where P is the projection $\mathfrak{H}' \rightarrow \mathfrak{H}$. The converse that T constructed in that manner is a POV-measure is also true.

Now Jauch and Piron made the following conjecture: For every generalized spectral measure $h: B(X) \rightarrow \mathcal{E}(\mathfrak{H})$ which satisfies (1)–(3) there exists an extension

$\mathfrak{S} \supset \mathfrak{H}$ and a spectral measure $h': B(X) \rightarrow \mathcal{E}(\mathfrak{S}')$ so that $h(\Delta) = P \wedge h'(\Delta)$ where P is the projection $\mathfrak{S}' \rightarrow \mathfrak{S}$.

By the theorem of Neumark there must be a POV-measure on $B(X)$.

The existence of this measure can be proved without the detour over h satisfying (1)–(3). It suffices to define the localizability of a system S in a space X in the following manner:

(A1) for every $\Delta \in B(X)$ and every state $s \in \mathcal{S}$ of the system S there is a probability $p(\Delta, s)$ for S in Δ . This function is additive in Δ and linear in s so that $\Delta_1 \cap \Delta_2 = \emptyset \Rightarrow p(\Delta_1 \cup \Delta_2, \cdot) = p(\Delta_1, \cdot) + p(\Delta_2, \cdot)$ and $s_1, s_2 \in \mathcal{S}, \alpha, \beta \geq 0, \alpha + \beta = 1 \Rightarrow p(\cdot, \alpha s_1 + \beta s_2) = \alpha p(\cdot, s_1) + \beta p(\cdot, s_2)$.

In order that this probability may be experimentally determined, we must suppose

(A2) for every Δ there is an observable $P(\Delta)$ with the mean value $p(\Delta, s)$ in every state s .

The definition requires that the sum of the observables $P(\Delta_1), P(\Delta_2)$ should exist, at least for $\Delta_1 \cap \Delta_2 = \emptyset$ (cf. [4]). In the sequel we shall deal with the case that \mathcal{E} is a Boolean algebra or the lattice of projections in a Hilbert space \mathfrak{S} , and so this condition will be satisfied.

In the second case the observables $P(\Delta)$ are operators on \mathfrak{S} and, in order that the mean value be from the interval $\langle 0, 1 \rangle$ for every state, it is necessary and sufficient that $P(\Delta)$ be a positive operator and $P(\Delta) \leq I$ (I – the unit projection). The condition $P(X) = I$ is equivalent to $p(X, s) = 1$ for every state s .

In the classical case, where \mathcal{E} is a Boolean algebra, we can consider the sets in a compact space T in place of \mathcal{E} (Theorem of Stone). $P(\Delta)$ are then measurable functions on T and states are measures on T . The necessary and sufficient condition that the mean value lie in $\langle 0, 1 \rangle$ is that the values of P lie in $\langle 0, 1 \rangle$. This follows from the fact that for every event $x \in \mathcal{E}$ it is possible to find a state $s \in \mathcal{S}$ that $s(x) = 1$ provided that the structure $(\mathcal{E}, \mathcal{S})$ is strongly-order – determining [3] (cf. [4]). In this case, therefore, the localizability is a system of functions $P(\Delta), \Delta \in B(X)$, with values in $\langle 0, 1 \rangle$ and $P(\Delta_1 \cup \Delta_2) = P(\Delta_1) + P(\Delta_2)$ for $\Delta_1 \cap \Delta_2 = \emptyset$. $P(X) = 1$ is always equivalent to S in X .

Now the following statement is valid (this statement is an analogue of the statement from [2]):

It is possible to find for \mathcal{E} an extension \mathcal{E}' and a system of states \mathcal{S}' such that there is an isomorphism $h: B(X) \rightarrow \mathcal{E}'$ and for every state $s \in \mathcal{S}$ there is a state $s' \in \mathcal{S}'$ so that the probability $S \in \Delta$ in the state s is equal to the probability of $h(\Delta)$ in the state s' .

Proof. We shall consider the representation of the algebra \mathcal{E} in the space T . In the product $X \times T$ we have sets $\Delta \times M, \Delta \in B(X), M \in \mathcal{E}$. Now we define the measure s' on $X \times T$ for $s \in \mathcal{S}$

$$s'(\Delta \times M) = \int_M P(\Delta) ds.$$

s' is a measure because P is additive in \mathcal{A} . We can extend s' on the algebra \mathcal{E}' generated by the sets of the form $\Delta \times M$. In this algebra \mathcal{E}' the sets $X \times M$ form an subalgebra that is isomorphic to \mathcal{E} . For every Δ we define $h(\Delta) = \Delta \times T$ and h is an isomorphism. Finally it is

$$s'(h(\Delta)) = \int_T P(\Delta) ds = p(\Delta, s).$$

This completes the proof.

The part about the states is valid for [2], too:

It is possible to extend every state s in \mathfrak{S} to a state s' in \mathfrak{S}' so that we have $s'(h(\Delta)) = s(P(\Delta))$ (in [2] there is F in place of h and B in place of P).

Proof. When the state s is irreducible then there is a vector $\varphi \in \mathfrak{H}$ that defines this state. $h(\Delta)$ is formed by all sums $\sum_k \Delta_k \cdot x_k$ (cf. [2]) where $\Delta_k \subset \Delta$. It follows from this that $(X - \Delta)\varphi$ is orthogonal to $h(\Delta)$ and from the equality $\varphi = X\varphi = \Delta\varphi + (X - \Delta)\varphi$ it follows that $\Delta\varphi$ is the projection of φ into $h(\Delta)$. From the definition of the scalar product in \mathfrak{H}' the equation

$$(h(\Delta)\varphi, \varphi)_{\mathfrak{H}'} = (P(\Delta)\varphi, \varphi)_{\mathfrak{H}}$$

follows. The validity for composed states is a consequence of linearity.

References

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- [3] *J. C. T. Pool: Baer*-Semigroups and the Logic of Quantum Mechanics, Comm. Math. Physics 9 (1968), 118.*
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Souhrn

O ZOBECNĚNÉ LOKALIZOVATELNOSTI

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Ukazuje se, že definice zobecněné lokalizovatelnosti podle Jaucha a Pirona může být přímo zdůvodněna. Je dokázána věta pro klasické systémy, která je analogem Neumarkovy věty o reprezentaci POV měř.

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